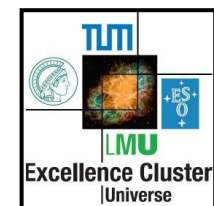




Correlations and clustering in nuclear matter and heavy ion collisions



Hermann Wolter, University of Munich, Germany



5th Int. Symposium on the Nuclear Symmetry Energy (NuSYM15)
Krakow, Poland , June 29 – July 2, 2015

Correlations and clustering in nuclear matter and heavy ion collisions

very wide subject: introductory remarks and overview,
subjects discussed in much more detail in later talks

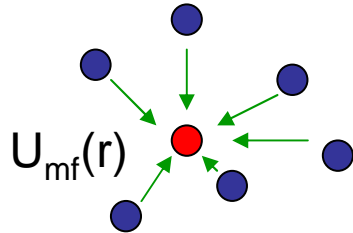
Aim of this talk:

- emphasize the importance of correlations and clustering in the study of the symmetry energy
- discuss what is involved in a proper treatment
- stress simple concepts (more detail in later talks)

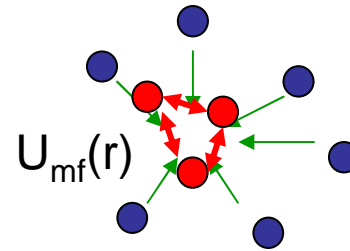
I will use the results and graphs of many workers in the field
(and apologize, if not always cited properly).

the two
main players:

mean field (mf) ↔ clustering



particle in a mean field,
one-body approach



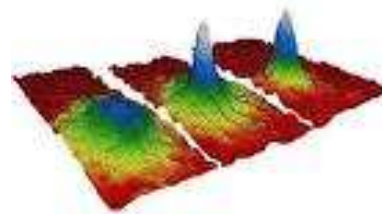
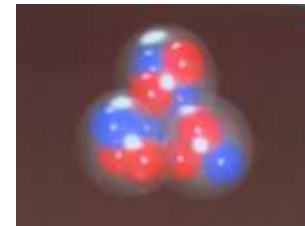
few-body correlation
in a mf
(medium-modification of cluster)

Examples:

1.) e.g. pairing. can be converted into a one-quasi-particle picture.
→ well studied

2.) quartetting, α -correlations

3.) BEC, other fields,...



... correlations: the seeds to clustering and cluster production

Clustering \longleftrightarrow Symmetry Energy. Relevance to NuSYM15??

nucleons come in two flavors: n, p

→ nn,pp interactions different from pn interaction (stronger)

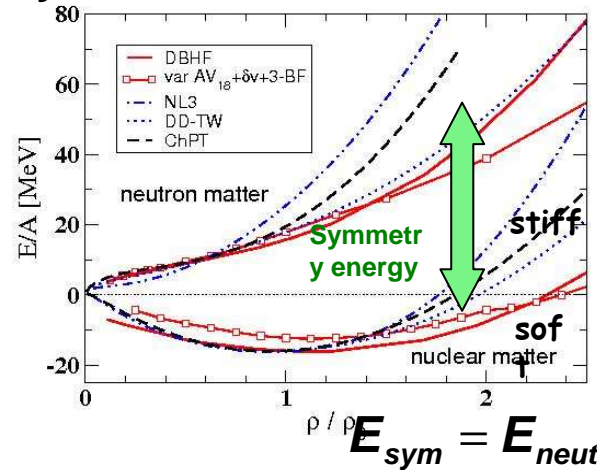
→ in asymmetric system: → U_p and U_n different → symmetry energy

density-
asymmetry dep.
of nucl.matt.

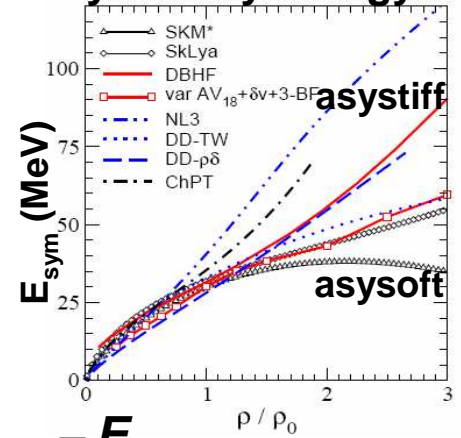
$$\varepsilon(\rho_B, \delta) = \varepsilon_{nm}(\rho_B) + E_{sym}(\rho_B) \delta^2 + O(\delta^4) + \dots \quad \delta = \frac{\rho_n - \rho_p}{\rho_n + \rho_p}$$

$$U(\rho, k; \delta) = \frac{\partial \varepsilon(\rho, \delta)}{\partial f(\rho, k)} = \underbrace{U_0(\rho, k)}_{U_\tau(\rho, k)} + \underbrace{U_{sym}(\rho, k)}_{U_\tau(\rho, k)} (\tau \delta) + \dots$$

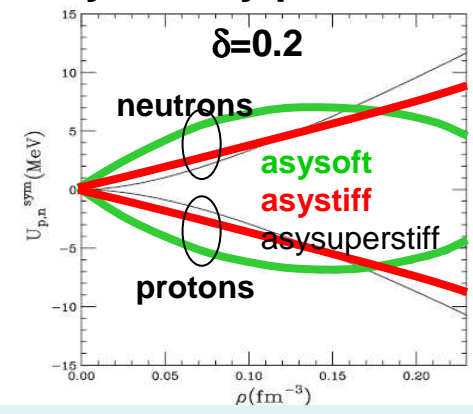
symmetric and neutron matter



symmetry energy



symmetry potential



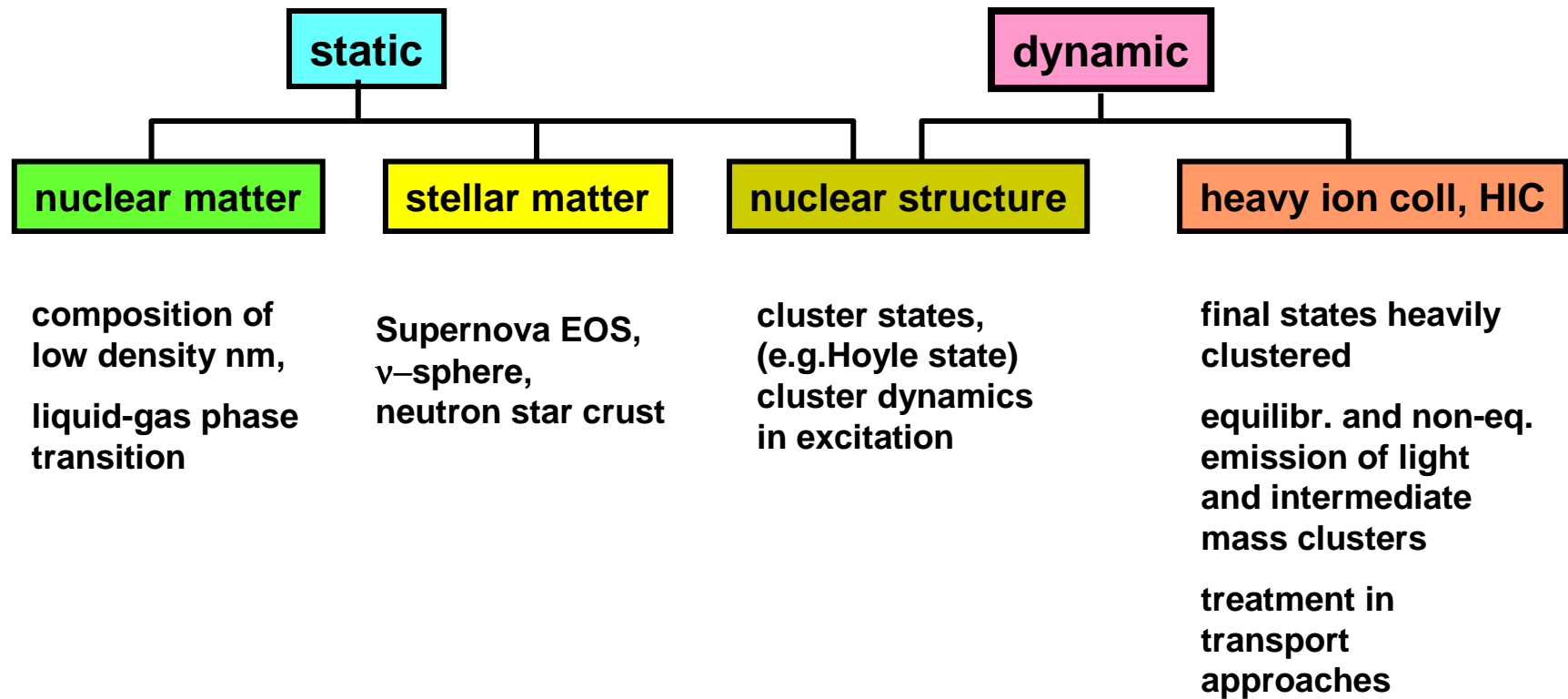
C. Fuchs, H.H. Wolter, EPJA 30(2006)5

..also momentum dependence
→ effective mass splitting

where do clusters come in?

1. clusters properties are driven by the symmetry energy, i.e. the N/Z ratio
2. isospin fractionation between clusters and gas
3. clusterization gives a direct contribution to the symmetry energy:
correlation depends on asymmetry of system; stronger in symmetric system

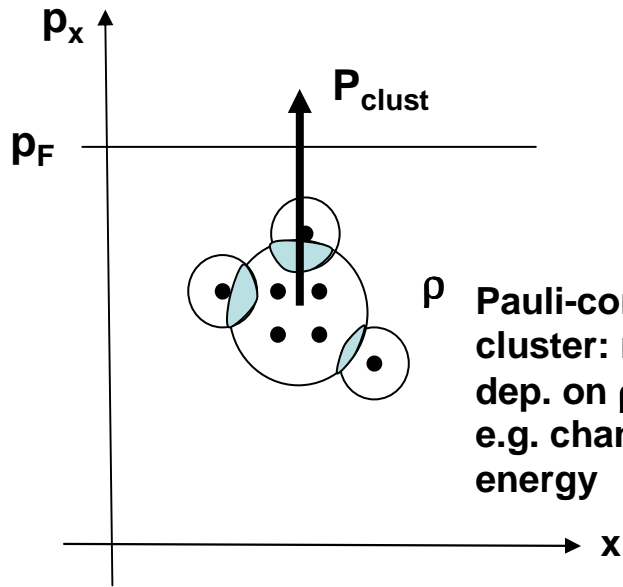
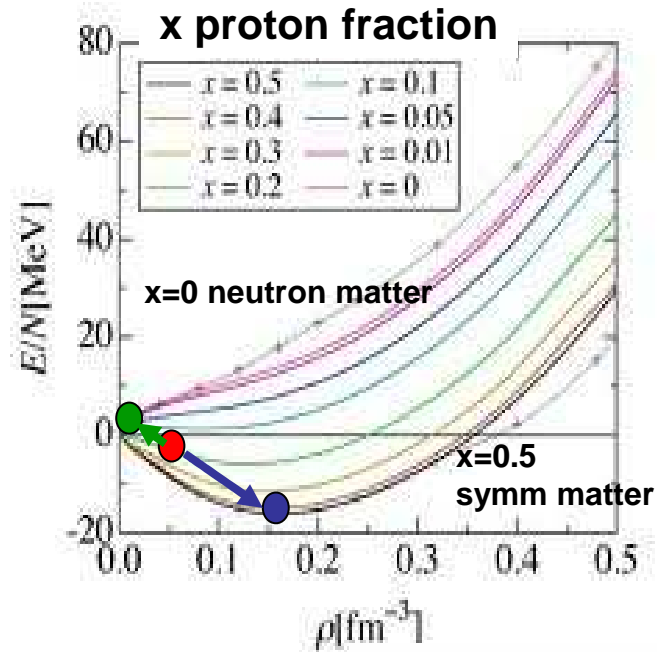
Attempt to Cluster Systematics in Nuclear Systems



→ This is a kind of outline of this talk

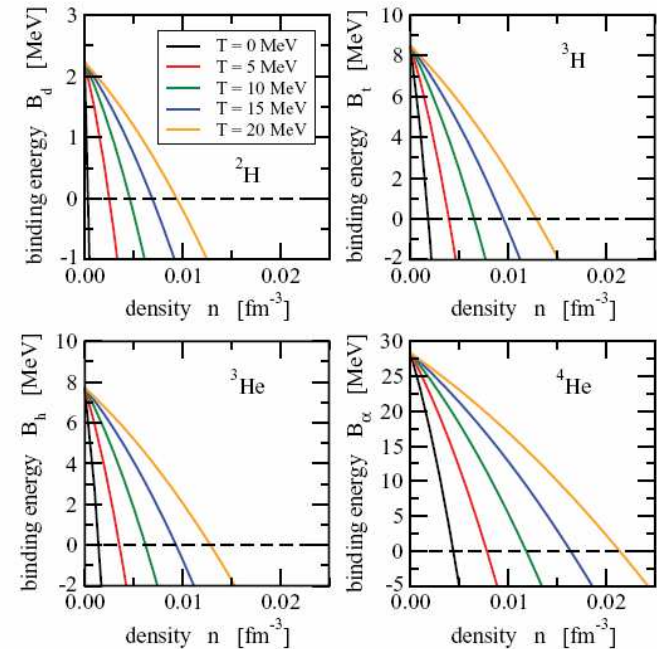
Composition of dilute matter

decrease energy by inhomogeneity
 → fractionation into clusters at
 higher density and neutron gas



Pauli-correlations weakens
 cluster: medium effects
 dep. on ρ , T , P_{clust} ,
 e.g. change of binding
 energy

clusterization increases with $\rho \searrow$, $T \nearrow$, $P_{cm} \nearrow$



1. Few-body Schrödinger eq. equation including blocking and in-medium (quasi-)particle energies (Röpke)

$$[E^{\text{qu}}(1) + \dots + E^{\text{qu}}(A) - E_{A,v}^{\text{qu}}(K)]\psi_{AvK}(1 \dots A) + \sum_{1' \dots A'} \sum_{i < j} [1 - \tilde{f}(i) - \tilde{f}(j)] V(ij, i' j') \prod_{k \neq i, j} \delta_{kk'} \psi_{AvK}(1' \dots A') = 0.$$

EoS in grand canonical picture: density as fct of chem. pot. and temperature

$$n_p^{\text{tot}}(T, \tilde{\mu}_p, \tilde{\mu}_n) = \frac{1}{\Omega} \sum_{A,v,K} Z f_{A,Z} [E_{A,v}^{\text{qu}}(K; T, \tilde{\mu}_p, \tilde{\mu}_n)], \quad \text{corresp. for neutrons}$$

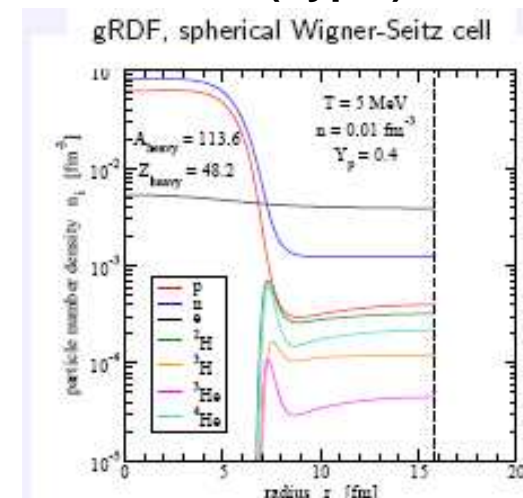
2. Nucl. statistical Equilibrium (NSE) (Hempel): energies of isolated A-nucleon clusters
 → incorrect high density limit, since clusters never disappear
 → remedy: excluded volume with increasing density. no space for clusters

3. Generalized rel mf (RMF) approach mit cluster degrees of freedom (Typel)

$$\mathcal{L}(\psi; \underbrace{\rho, \omega, \rho, \delta}_{\text{meson with DD coupling}}; \underbrace{\mathbf{d}, \alpha, t, {}^3\text{He}}_{\text{light cluster with DD masses}})$$

meson with DD coupling light cluster with DD masses
 interpolate correctly from zero to high density

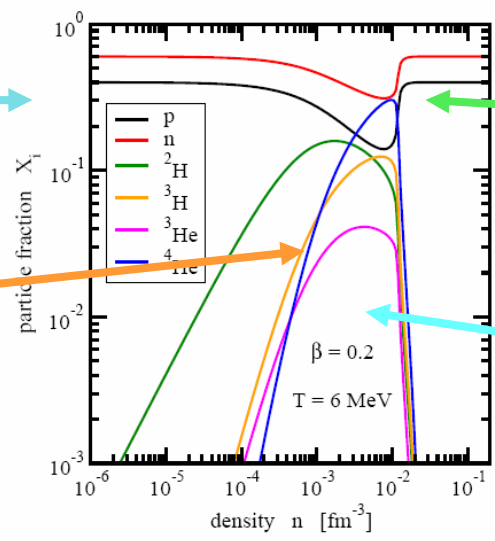
treatment difficult for heavier clusters:
 → Thomas-Fermi calculation in Wigner-Seitz cell
 (periodic space)



Particle Fractions

very low density: p,n

Increasing density:
clusters arise: deuteron
first, but then α
dominates



Mott density:
clusters melt,
homogeneous p,n
matter;

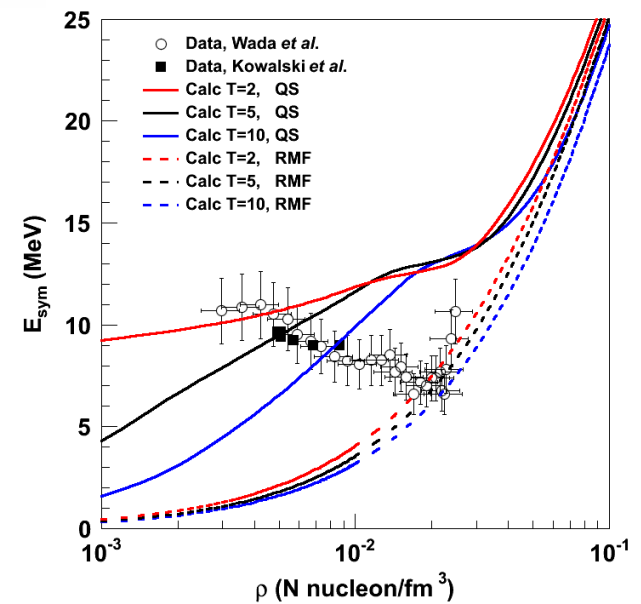
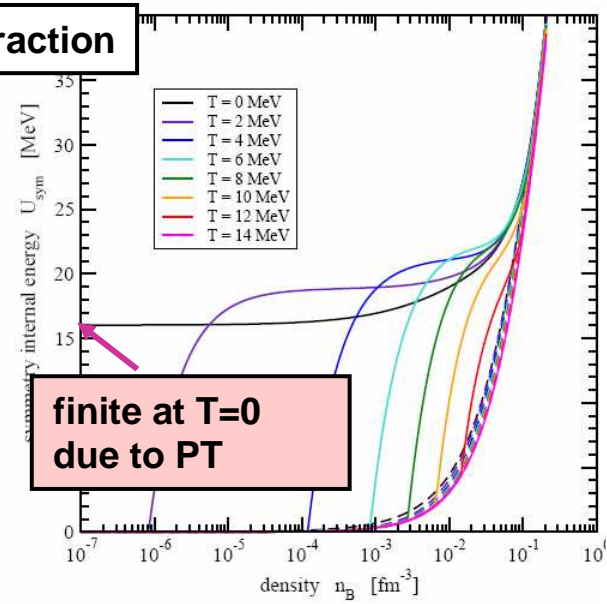
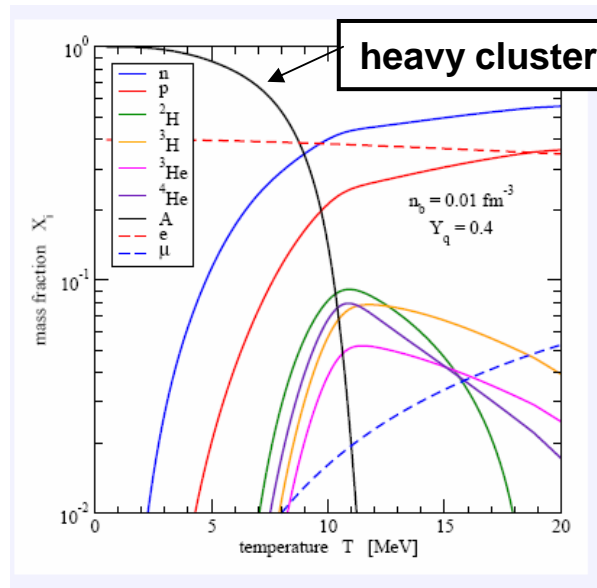
here heavier nuclei
(embedded into a
gas) become
important here

S.Tygel, G. Röpke, et al., PRC 81 (2010)

composition as fct of Temp.

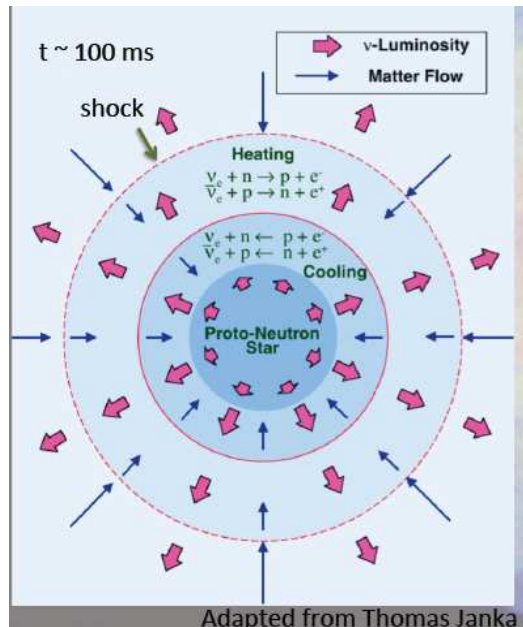
Symmetry Energy

comparison to data from
heavy ion collisions



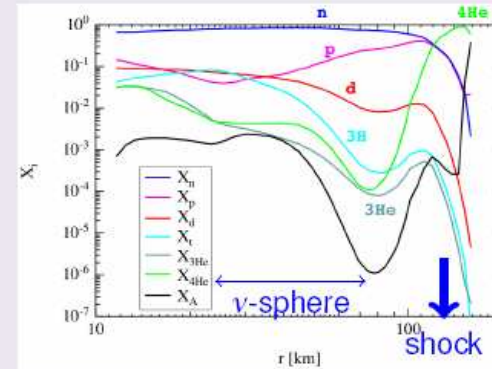
J. Natowitz, G. Röpke, ... HHW, PRL 104, 202501 (2010)

Core Collapse Supernovae and the Neutrinosphere



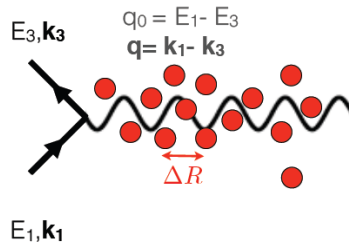
Mass fraction of light clusters in the post-bounce supernova core, based on nuclear statistical equilibrium.

Sumiyoshi and Röpke, PRC77 (2008) 055804.



Correlations & Neutrino Scattering

- Neutrinos “see” more than one particle in the medium.
- Nature of spatial and temporal correlations between nuclei, nucleons and electrons affect the scattering rate.



At small q_0 and q the neutrino cannot resolve single particles.

Sawyer (1975, 1989)
Iwamoto & Pethick (1982)
Horowitz & Vherberger (1991)
Raffelt & Seckel (1995)

Differential Scattering/Absorption Rate:

$$\frac{d\Gamma(E_1)}{d\cos\theta dq_0} = \frac{G_F^2}{4\pi^2} (E_1 - q_0)^2 [(1 + \cos\theta) S_V^{RPA}(q_0, q) + (3 - \cos\theta) S_A^{RPA}(q_0, q)]$$

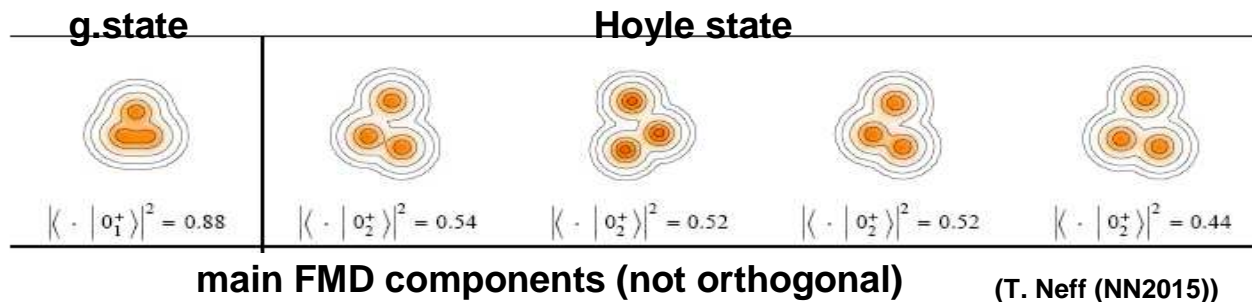
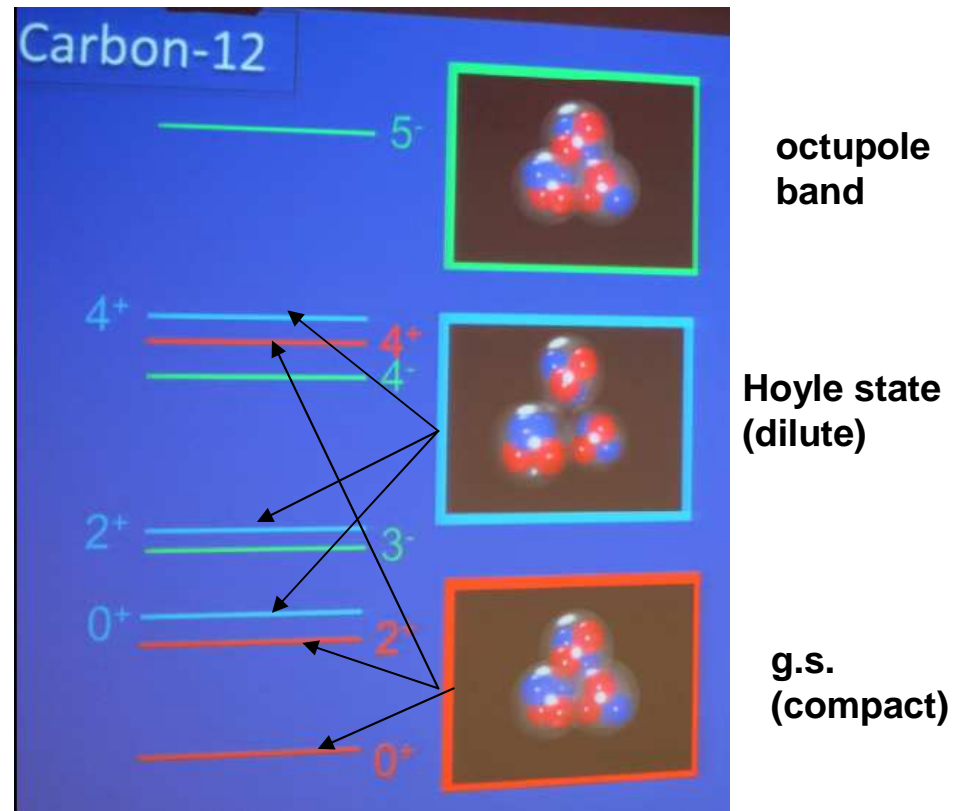
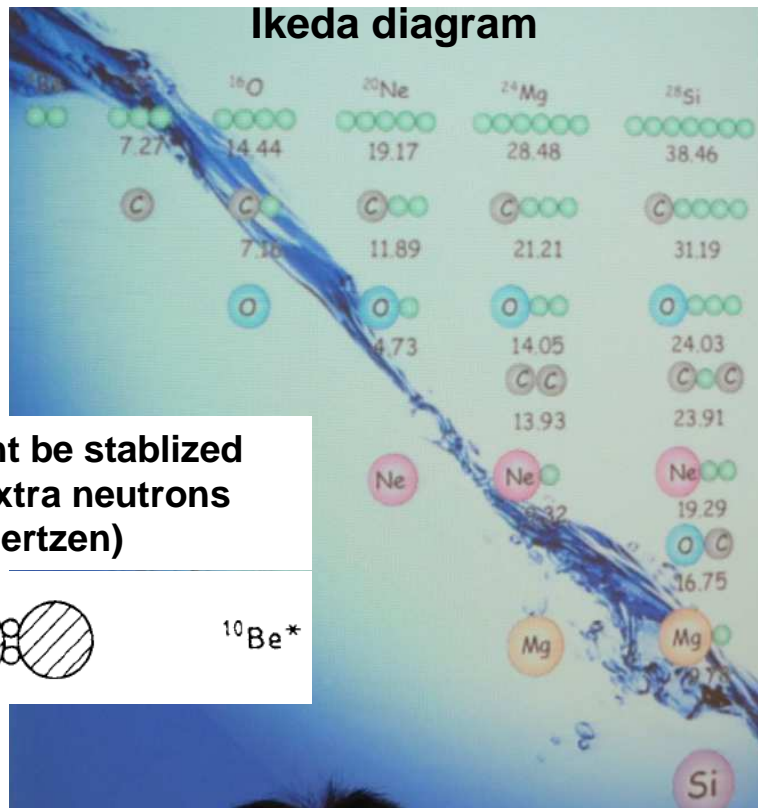
spectrum of density fluctuations

spectrum of spin fluctuations

direct interface with heavy ion physics, simulate in a box calculation!

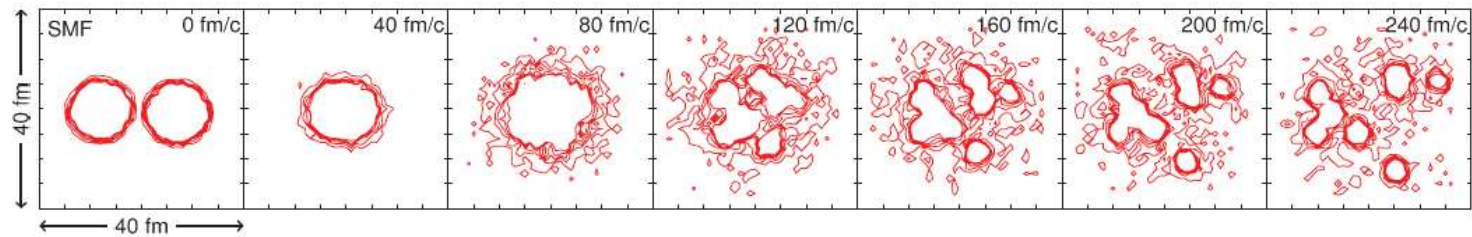
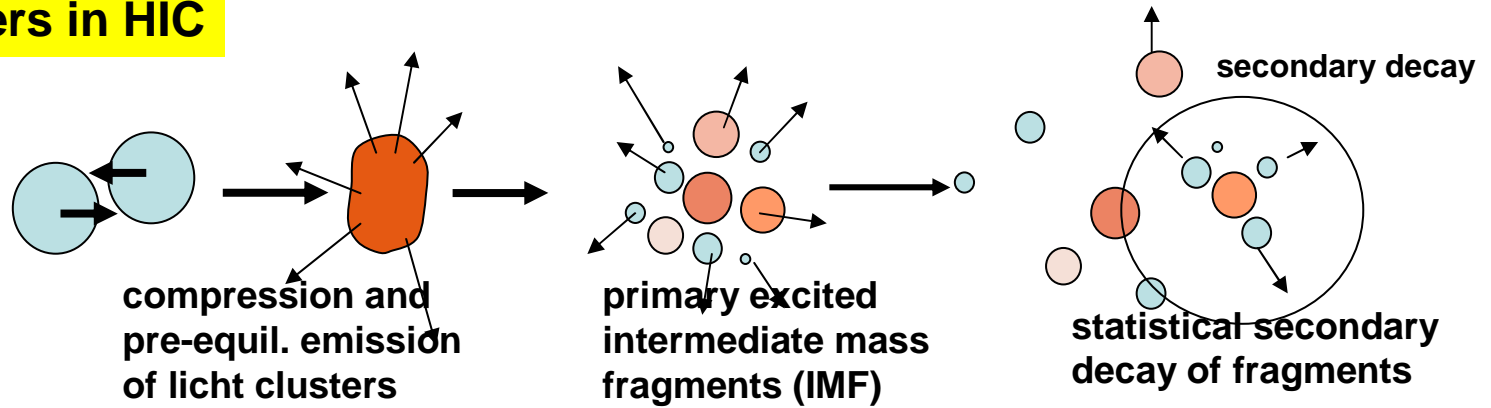
supernova ↔ femtonova

Clusters in nuclear structure



surface clustering (\rightarrow Typel): change of relation neutron skin \leftrightarrow slope of symmetry energy

Clusters in HIC

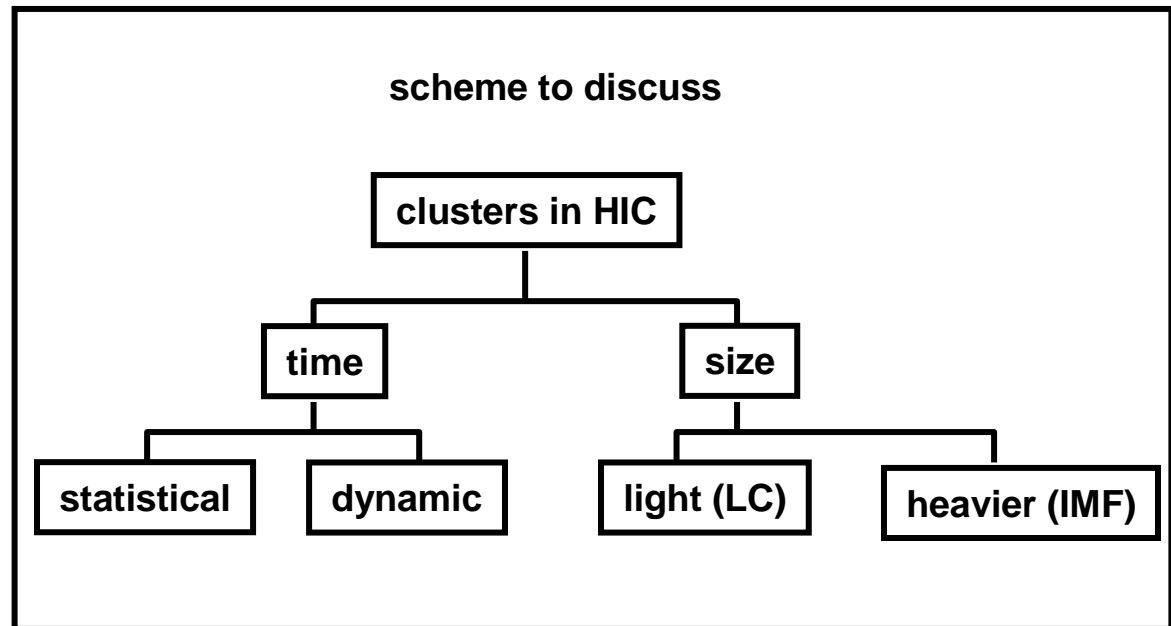


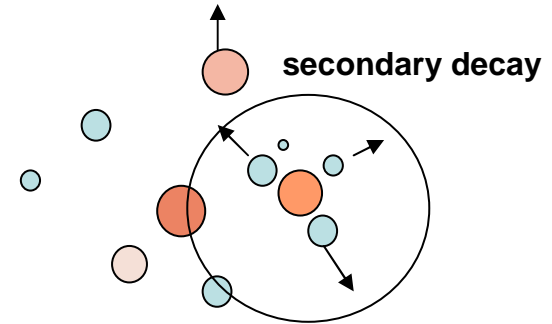
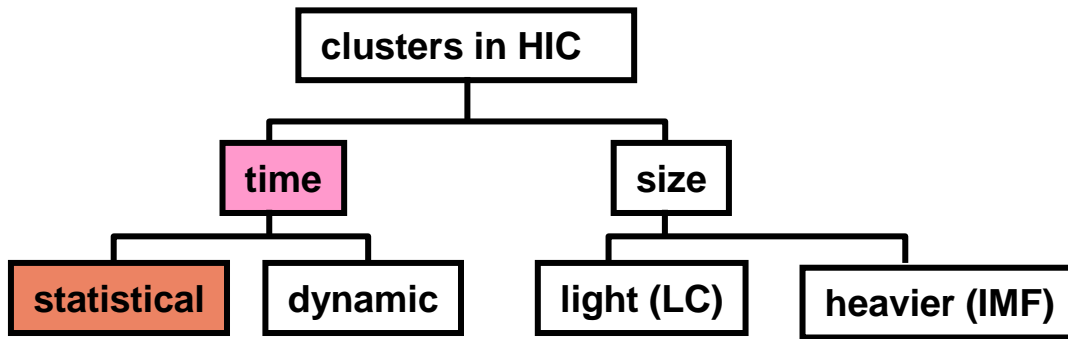
large fractions of particles in clusters, e.g.

	Partitioning of protons	
	Xe + Sn 50 MeV/u	Au + Au 250 MeV/u
p	≈10%	21%
α	≈20%	20%
d, t, ^3He	≈10%	40%
$A > 4$	≈60%	18%

INDRA data, Hudan et al., PRC67 (2003) 064613.

FOPi data, Reisdorf et al., NPA 848 (2010) 366.

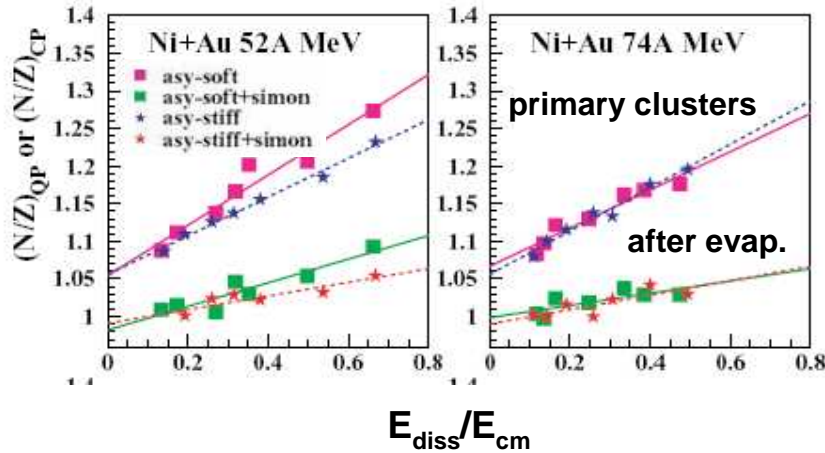




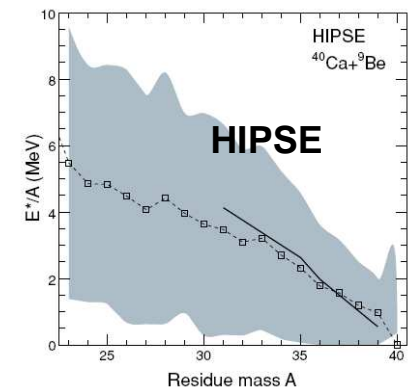
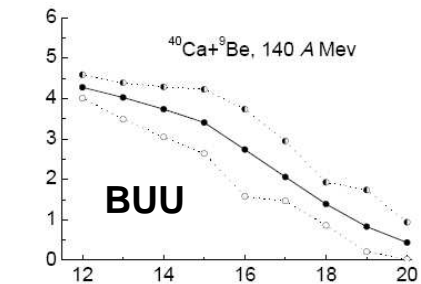
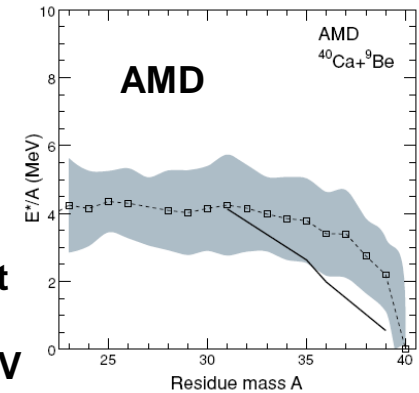
Freeze-out approximation: assume: A, Z, volume, excitation energy/A
statistical decay code: SMM, Gemini++, SIMON, etc
 „afterburner“

excitation energy
calculated in different
approaches for
 $^{40}\text{Ca}+^9\text{Be}$, at 140 AMeV

large influence on observables!
hope that „isospin blind“?

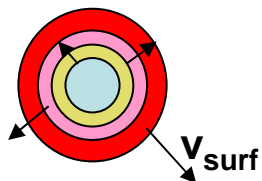
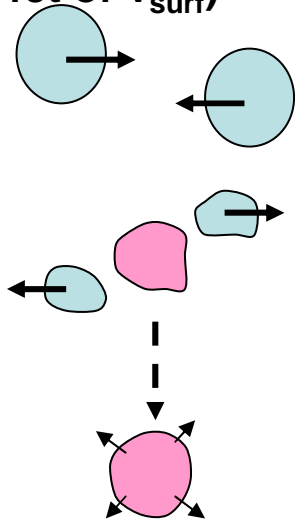


Galichet et al., Phys. Rev. C79, 064615 (2009)



Freeze-out and the Supernova (SN) ν -sphere in HI collisions?

Coalescence analysis (as fct of v_{surf})



„differential“ freeze-out analysis:

- source reconstruction,
- analysis as fct of $v_{surf} \sim$ time of emission
- determination of thermodyn. properties and symmetry energy as fct of v_{surf}

Experiment and interpret. $^{64}\text{Zn} + (^{92}\text{Mo}, ^{197}\text{Au})$ at 35 A MeV
 S. Kowalski, J. Natowitz, et al., PRC75 014601 (2007)
 J. Natowitz, G. Röpke, S. Typel, et al., PRL 104, 202501 (2010),
 R. Wada, et al., PRC 85, 064618 (2013)

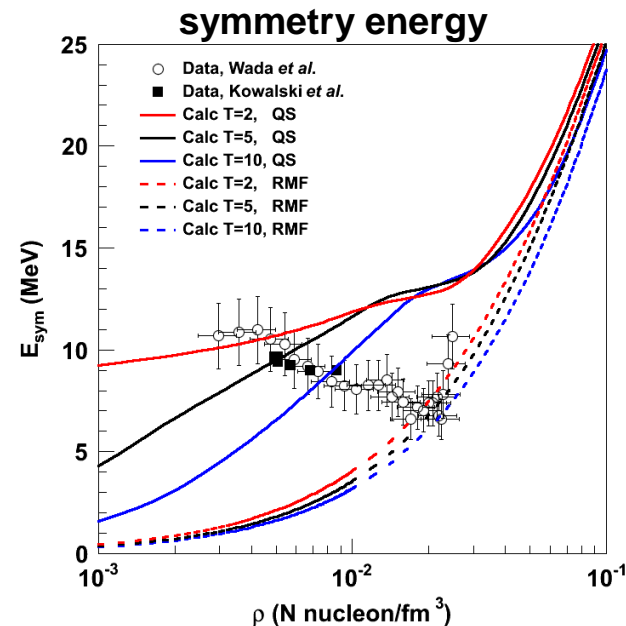
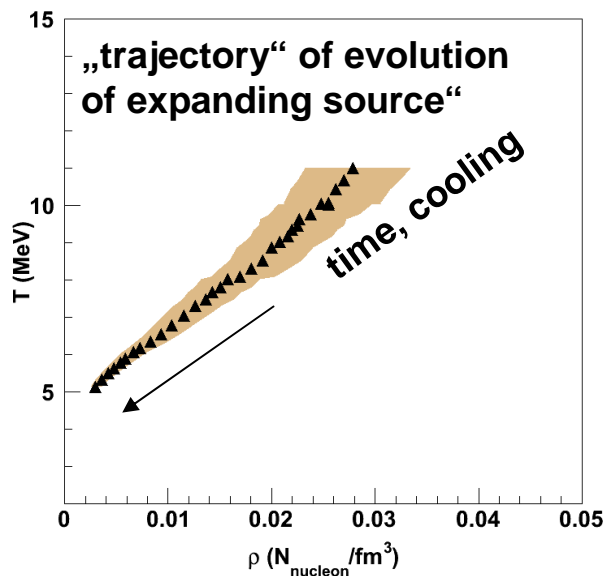
albergo thermometer

$$T_{HHe} = \frac{14.3}{\ln\left(\sqrt{9/8} \frac{1.59 Y(^2\text{H})Y(^4\text{He})}{Y(^3\text{H})Y(^3\text{He})}\right)}$$

density: momentum space coalescence model (Mekjian, PRC 17 (78))

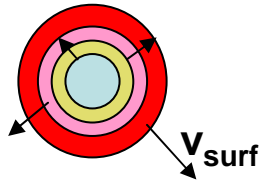
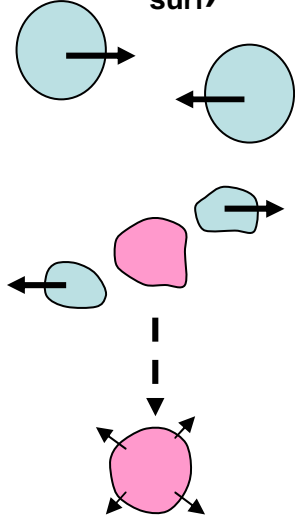
Symmetry free energy from isoscaling

$$\begin{aligned} \frac{Y_2}{Y_1} &= C e^{[\mu_2(n) - \mu_1(n)]N + [\mu_2(p) - \mu_1(p)]Z} / T \\ &= C e^{\alpha N + \beta Z} \\ \alpha &= 4F_{\text{sym}}[(Z_1/A_1)^2 - (Z_2/A_2)^2] / T \\ \beta &= 4F_{\text{sym}}[(N_1/A_1)^2 - (N_2/A_2)^2] / T, \end{aligned}$$



Freeze-out and the Supernova (SN) ν -sphere in HI collisions?

Coalescence analysis
(as fct of v_{surf})



„differential“ freeze-out analysis:
→ source reconstruction,
→ analysis as fct of

conditions of neutrinosphere:
densities $1/1000$ to $1/10 \rho_0$
temperature $T=4-5$ MeV
asymmetry $Y_e=0.1 - 0.25$

Experiment and interpret. $^{64}\text{Zn}+(^{92}\text{Mo},^{197}\text{Au})$ at 35 A MeV
S. Kowalski, J. Natowitz, et al., PRC75 014601 (2007)
J. Natowitz, G. Röpke, S. Typel, et al., PRL 104, 202501 (2010),
R. Wada, et al., PRC 85, 064618 (2013)

albergo thermometer

$$T_{HHe} = \frac{14.3}{\ln\left(\sqrt{9/8} (1.59) \frac{Y(^2\text{H})Y(^4\text{He})}{Y(^3\text{H})Y(^3\text{He})}\right)}$$

density: momentum space coalescence
model (Mekjian, PRC 17 (78))

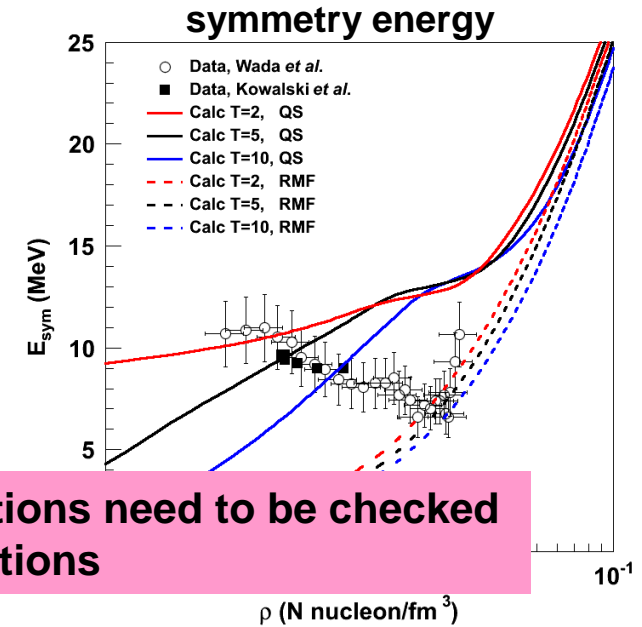
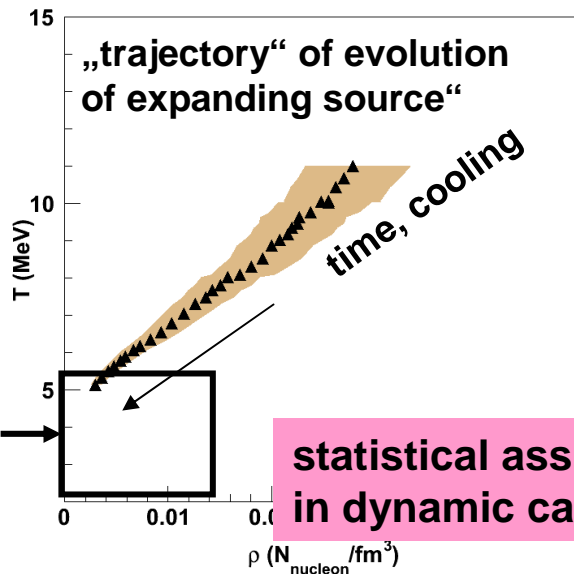
Symmetry free energy
from isoscaling

$$\frac{Y_2}{Y_1} = C e^{\{[\mu_2(n)-\mu_1(n)]N + [\mu_2(p)-\mu_1(p)]Z\}/T}$$

$$= C e^{\alpha N + \beta Z}$$

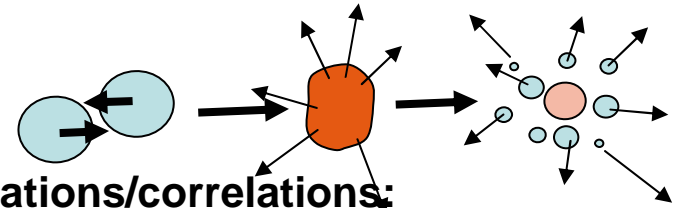
$$\alpha = 4F_{\text{sym}}[(Z_1/A_1)^2 - (Z_2/A_2)^2]/T$$

$$\beta = 4F_{\text{sym}}[(N_1/A_1)^2 - (N_2/A_2)^2]/T$$



statistical assumptions need to be checked
in dynamic calculations

Dynamical cluster formation in HIC



„dynamical clusters“: transport approach with fluctuations/correlations;
seeds of fragment/light cluster formation

BUU/BLE/BLOB

$$\left(\frac{\partial}{\partial t} + \frac{\vec{p}}{m} \vec{\nabla}^{(r)} - \vec{\nabla} U(r) \vec{\nabla}^{(p)} \right) f(\vec{r}, \vec{p}; t) = I_{coll} [\sigma^{in-med}] + \delta I_{fluct}$$

deterministic, dissipative 1-body equation + fluctuation

QMD/AMD

$$|\Phi\rangle = \mathcal{A} \prod_{i=1}^A \varphi(r; r_i, p_i) |0\rangle$$

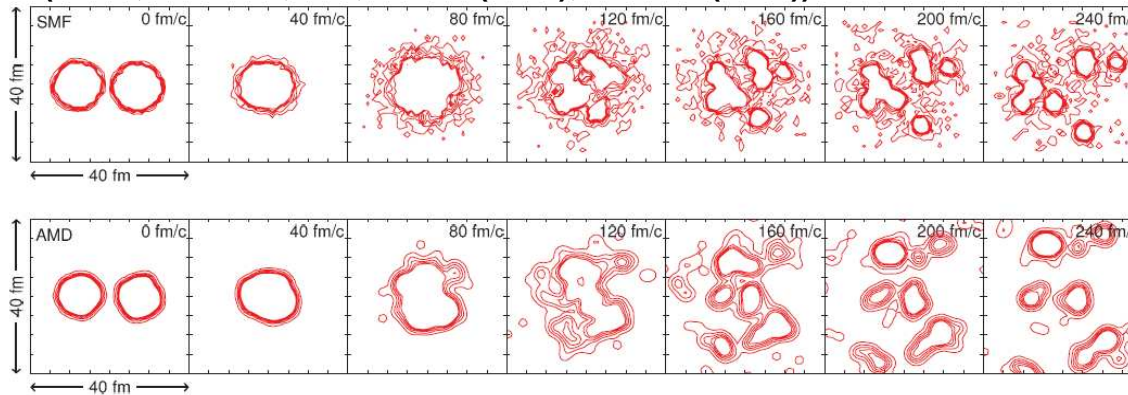
$$\dot{r}_i = \{r_i, H\}; \quad \dot{p}_i = \{p_i, H\}; \quad H = \sum_i t_i + \sum_{i,j} V(r_i - r_j)$$

TDHF + stochastic NN collisions

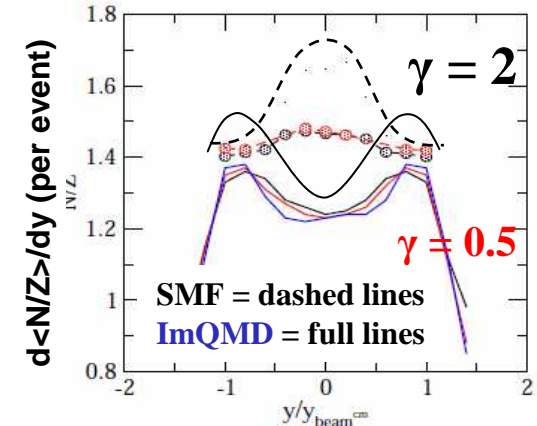
difference in spectrum of fluctuations

→ can be adjusted for IMF (intermediate mass fragments) formation, since IMF formation stabilized by mean field

Comparison of simulations: BUU(SMF)-AMD:
(Rizzo, Colonna, Ono, PRC76(2007); PRC82 (2010))



Comparison, SMF-ImQMD:
more transparency in QMD
(M. Colonna, X.Y.Zhang)



Interlude:

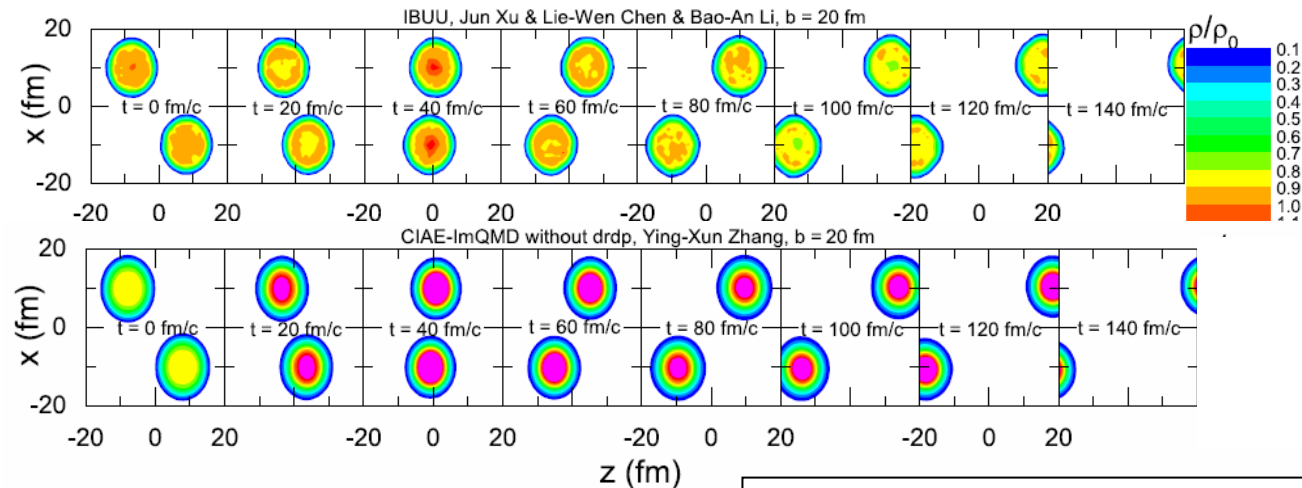
Code Comparison Project: Trento, ECT*, 2006 and 2009 Shanghai, Jan. 2014, Lanzhou 2014

check consistency of transport codes in calculations with same system (Au+Au), $E=100,400$ AMeV, identical (simple) physical input (mean field (EOS) and cross sections)

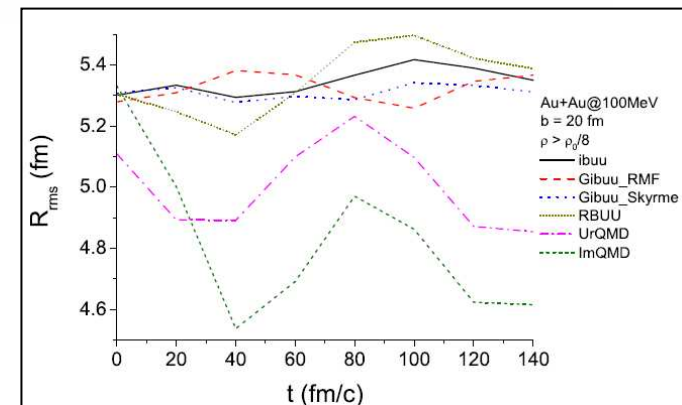
idea: establish sort of **theoretical systematic error or transport calculations** (and hopefully to reduce it)

1. step: Initialize colliding nuclei. usually not exact ground states

free
propagation
(large
impact
parameter)



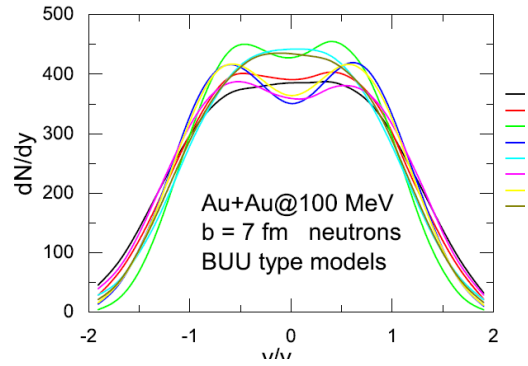
- nuclei oscillate, influences dynamical evolution in collision, part. at lower incident energies
- construct better gs, e.g. Thomas-Fermi



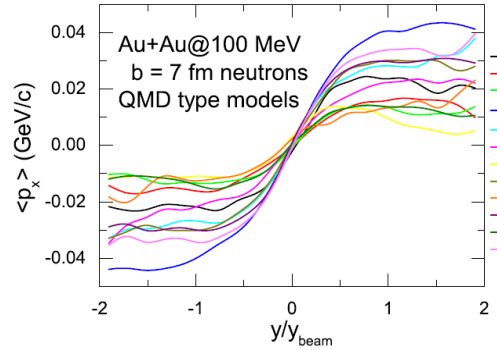
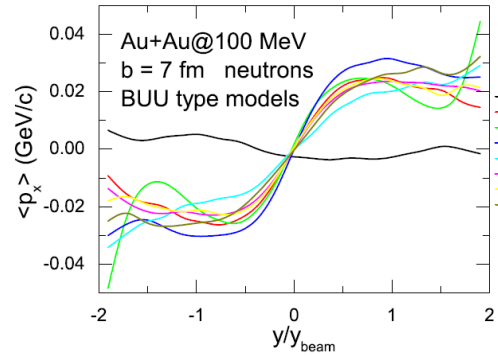
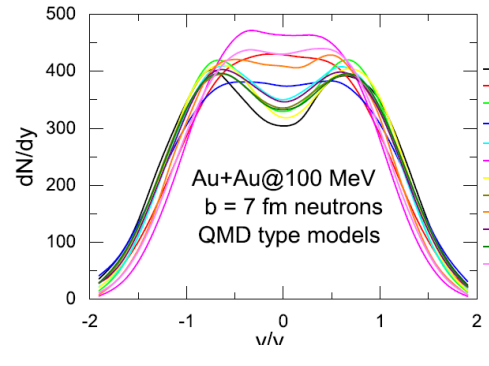
Examples of results: Au+Au, **PRELIMINARY**

E/A=100 MeV

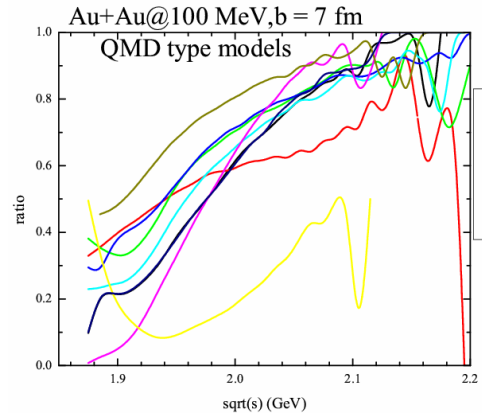
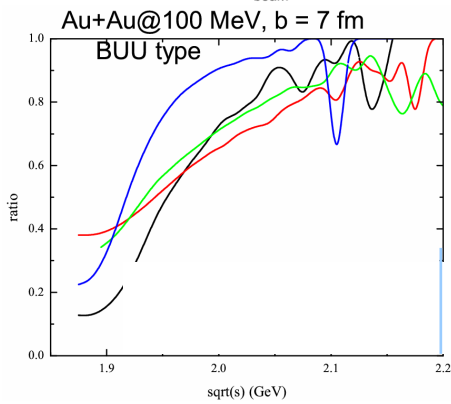
BUU models



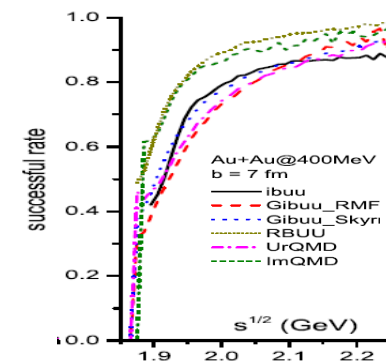
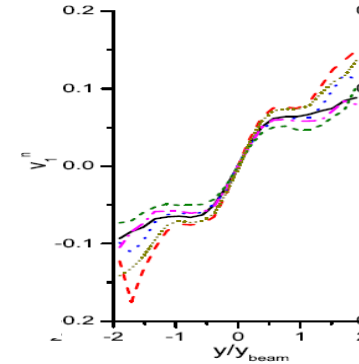
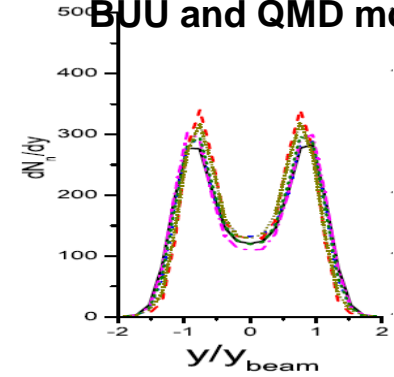
QMD models



successful coll. rate

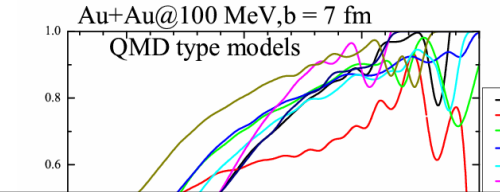
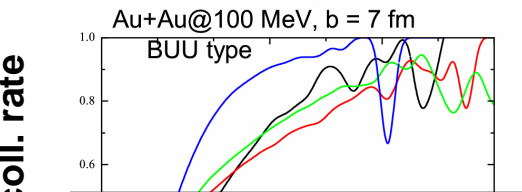
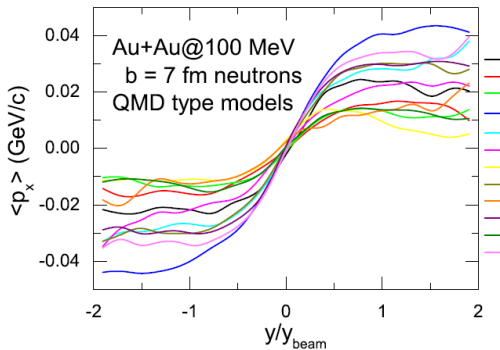
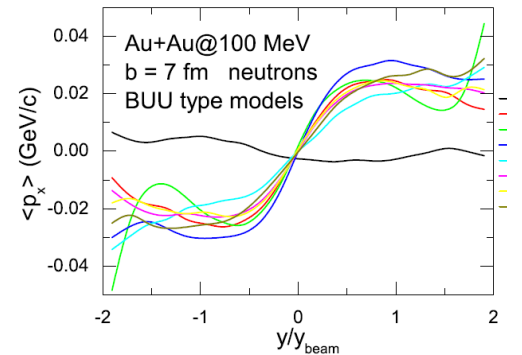
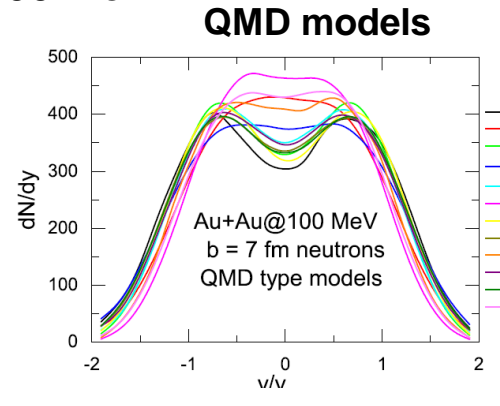
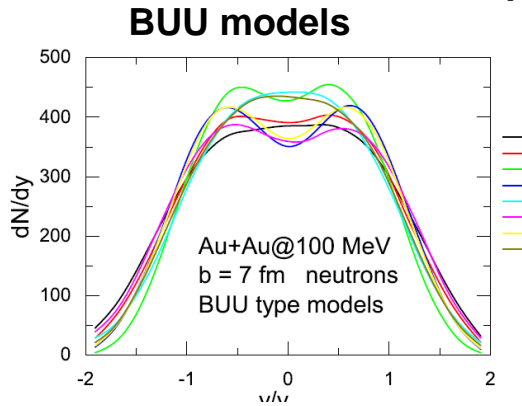


E/A=400 MeV
BUU and QMD models

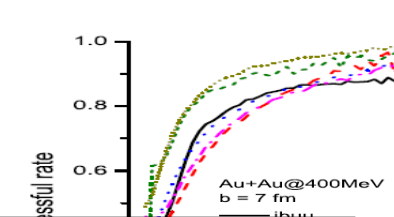
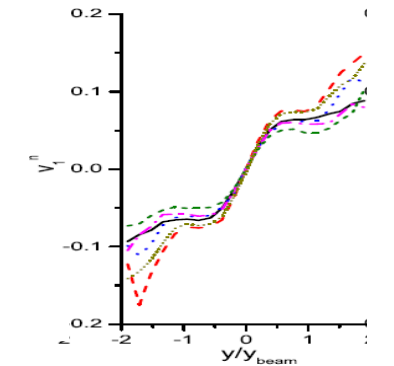
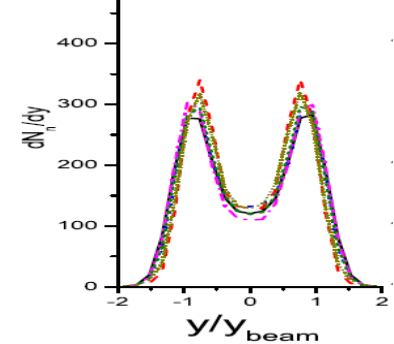


Examples of results: Au+Au, **PRELIMINARY**

E/A=100 MeV



E/A=400 MeV
BUU and QMD models



successful coll. rate

- considerable differences
- partly due to initialization, **but most likely mainly to collision term**
- no essential difference between BUU and QMD models
- 100 MeV sensitive region for flow because of competition between mf and collisions, comparison better at higher energy

End of Interlude on code comparison

MF
kyn
2.2

Cluster recognition algorithms:

MD: phase space connection: MST, SACA

BUU: a) coalescence: find contours of density $\rho_c \sim 1/10 \rho_0$

b) Test particle distribution sampling:

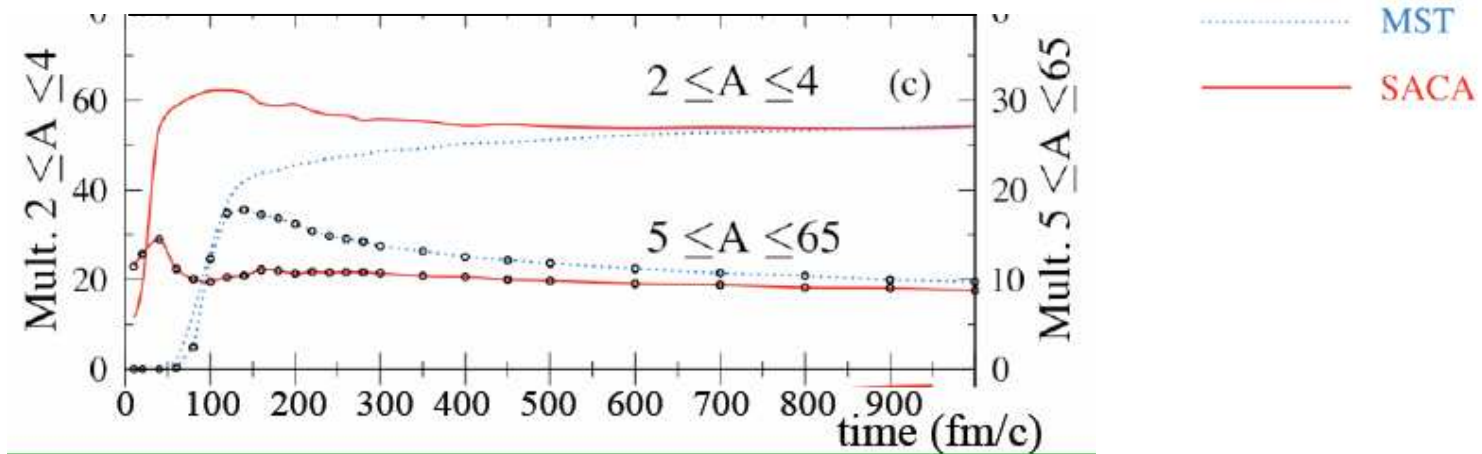
choose A out of $N_{TP} * A$ test particles with correct global properties.

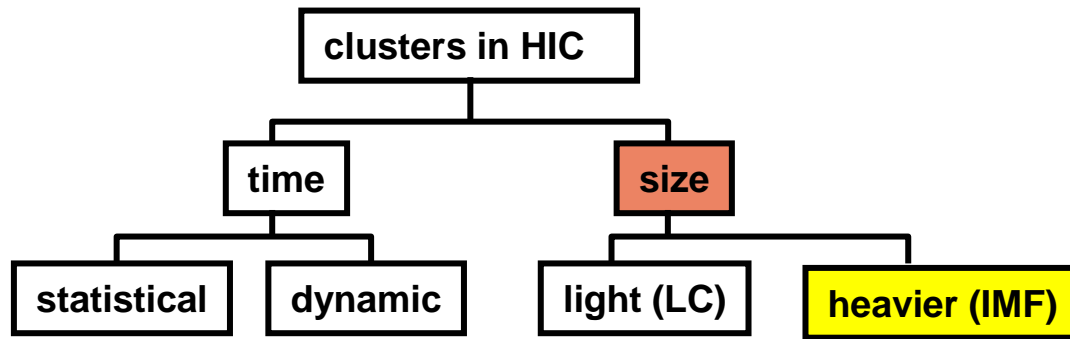
Treat these as nucleons and do coalescence or MST algorithm or SACA (as in QMD) and do this many times (~ 1000) to generate a distribution

→ similar in both approaches, but a-posteriori

→ not the same as dynamical clusters

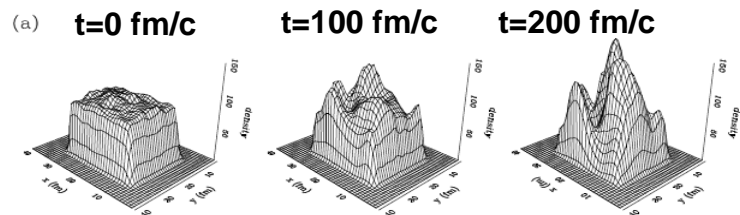
clusters identified earlier, but converge asymptotically and do not influence dynamics





IMF's: formation dominated by mean field, which favors matter at normal density

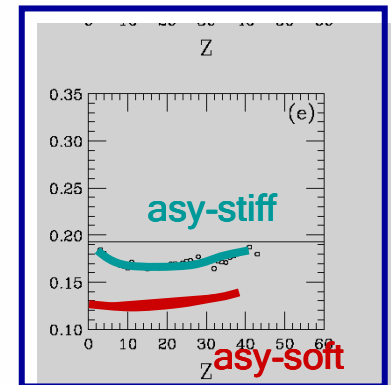
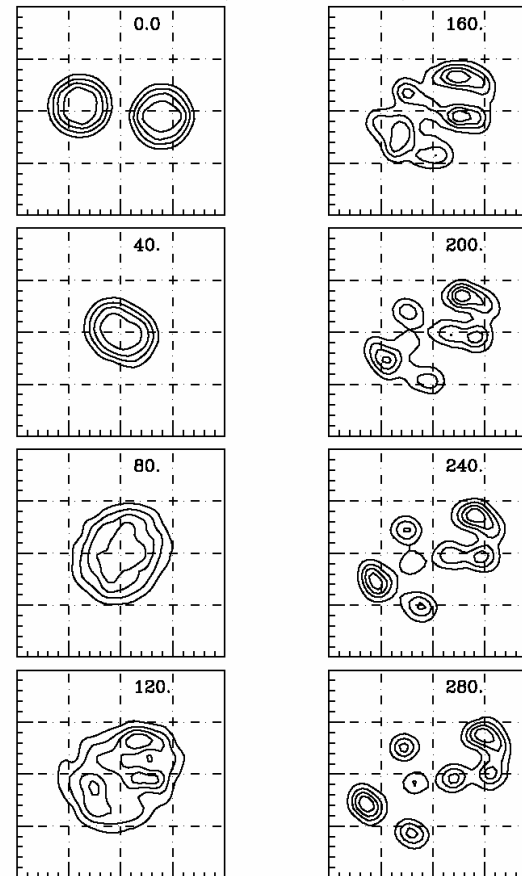
BUU calculation in a box (i.e. periodic boundary conditions) with initial conditions inside the instability region: $\rho = \rho_0/3$, $T = 5$ MeV, $\delta = 0$



→ Formation of „clusters (fragments)“, from small (numerical) fluctuations in the density. Time scale = growth time of the instable modes (V. Baran, et al., Phys.Rep.410,335(05))

successful applications for several observables: isospin transport and diffusion, liquid-gas phase transition, etc.

$^{124}\text{Sn} + ^{124}\text{Sn}$, 50 A MeV, central



Multifragmentation: Isospin fractionation at low densities

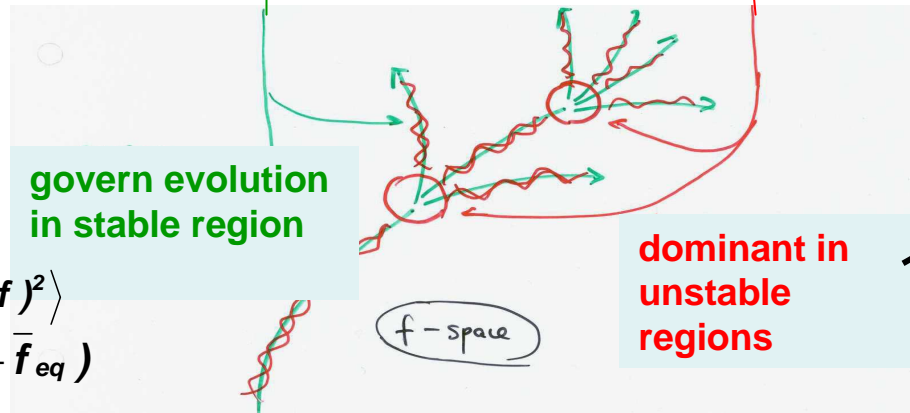
Treatment of Fluctuations (esp. in BUU)

$$f(r, p, t) = \bar{f}(r, p, t) + \delta f(r, p, t)$$

Mean field evolution (dissipative) Fluctuations (higher order correlations)

Boltzmann-Langevin eqn.

$$\frac{df}{dt} = I_{coll} + I_{fluc}$$



$$f = \bar{f} + \delta f; \quad \sigma^2 \equiv \langle (\delta f)^2 \rangle$$

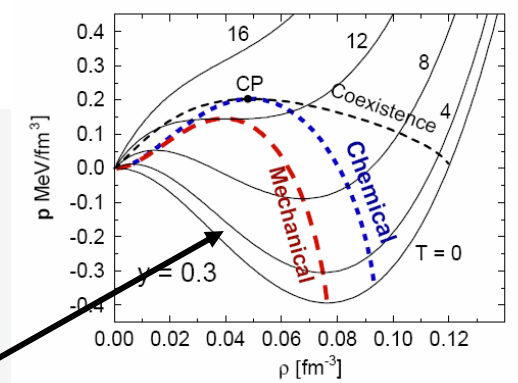
$$\sigma^2 \rightarrow \sigma_{eq}^2 = \bar{f}_{eq}(1 - \bar{f}_{eq})$$

approaches:

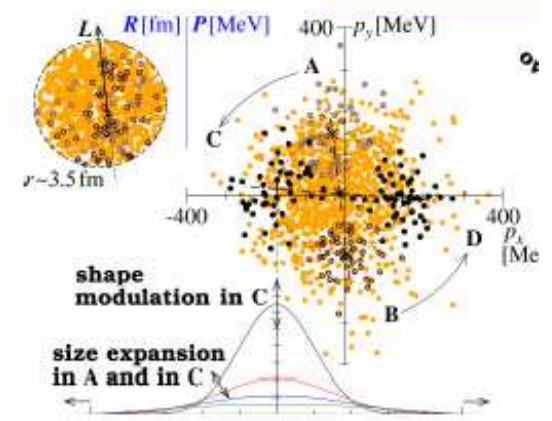
BOB (One-Body-Brownian): replace fluctuation term by fluctuating force, gauged to most unstable mode: Colonna, Guarnera

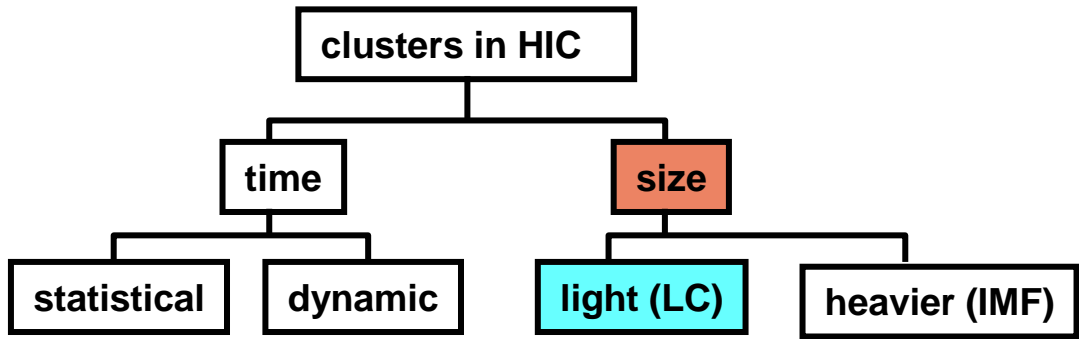
Stochastic MF dynamics (SMF): introduce locally statistical fluctuations into the phase space distribution at certain times according to $\sigma^2=f(1-f)$ projected on density fluctuations: Colonna, DiToro, HW

BLOB (Boltzmann-Langevin One-Body)
 Bertsch method, developed by M. Colonna, P. Napolitano (see talk later), fluctuation in full phase space
 One TP collision moves N_{TP} TP nearby in phase space, to simulate collisions of nucleons (P. Napolitano, this meeting).



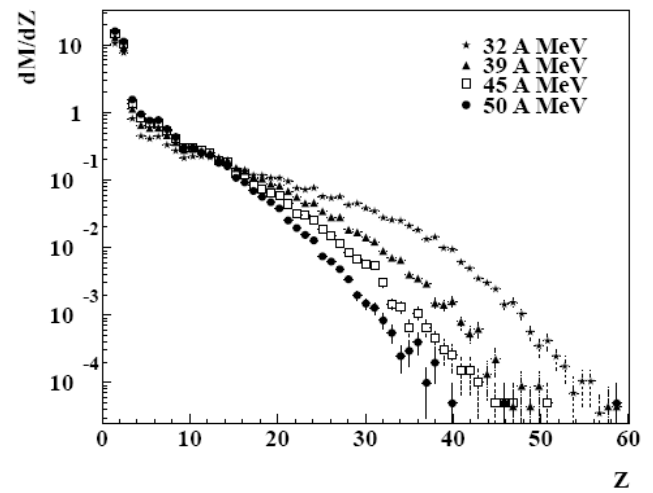
At a given time t ,
 in (r_a, p_a) ,
 for elastic coll. : $\dot{f}_a(r_a, p_a) = g$





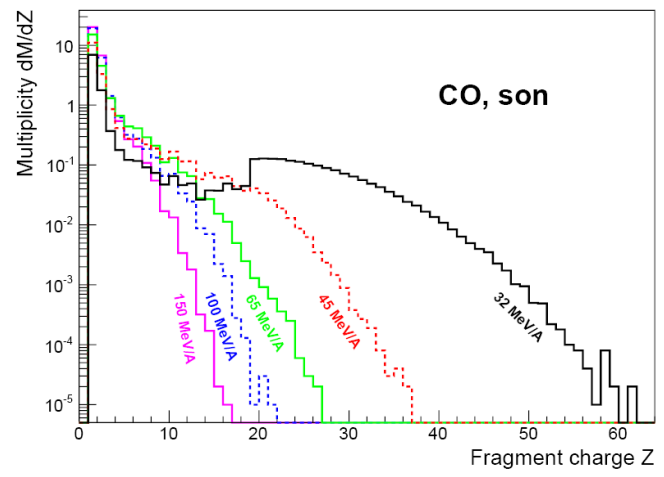
indication: overprediction of :nucleons and light clusters

$^{136}\text{Xe} + ^{124}\text{Sn}$, $E = 32, \dots, 150 \text{ A MeV}$



S. Hudan, et al., PRC67, 064613

$^{136}\text{Xe} + ^{124}\text{Sn}$, $b = 2 \text{ fm}$



- LC are correlation dominated (esp. Pauli-correlation).
- Not well described in BUU and MD models, since simple interactions and classical phase space distribution give bad eigenstates for LC
- need special treatment

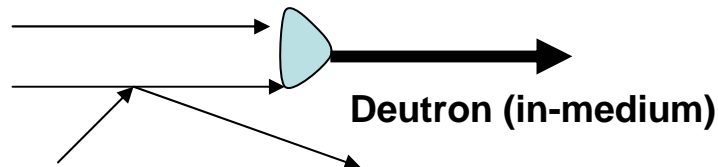
Treatment of Light Cluster dynamics in HIC:

circumvent: compare to „coalescent invariant“ cross sections
only justified if clusters play no dynamical role

solutions different in BUU and MD models:

Solution for BUU models:

LC distribution functions as explicit
degrees of freedom of type $NNN \rightarrow N\Delta$
(P. Danielewicz and Q. Pan, PRC 46 (1992))
(d,t,3He, **but no a!**)
→ coupled transport equations



Caveat: Medium properties of LC:

Medium corrections in the formation of light charged particles in heavy ion reactions

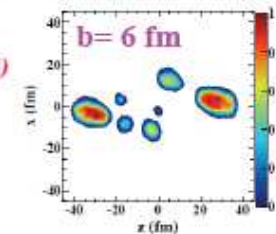
[C. Kuhrts](#), Beyer, Danielewicz,..PRC63 (2001) 034605

Calculated in nuclear matter and static nuclei in Generalized RMF approach by Typel, Röpke, et al., PRC81 (2010)

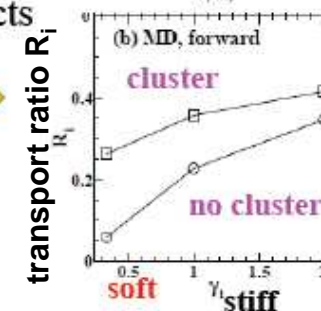
similar: transitions amplitudes in medium

pBUU calculations (Danielewicz)

*Coupland et al.,
PRC 84, 054603 (2011)*



Cluster effects
(d,t,He³)



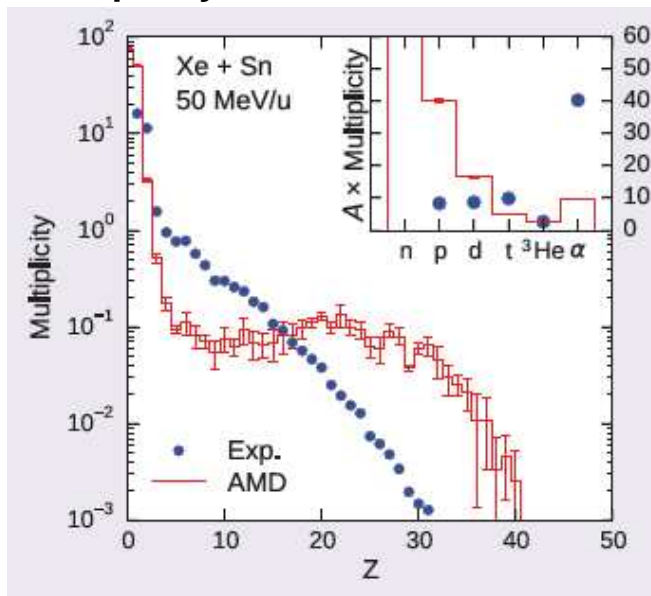
R_i : isospin transport ratio for charge equilibration in HIC between nuclei with different isospin content
e.g. $^{112,124}\text{Sn} + ^{112,124}\text{Sn}$
(MSU experiment)

Treatment of Light Cluster dynamics in HIC:

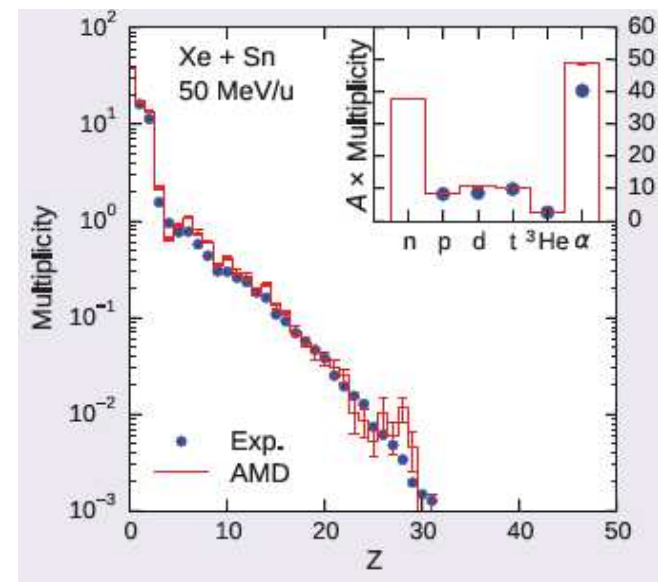
Solution für **AMD** (see talk of Ono)

1. in collision term consider formation of clusters in terms of overlap with cluster wave function (detailed balance?)
2. Manipulate phase space: put wave packets of cluster constituent in one place (conservation laws?)
3. consider Pauli correlation fully
4. include also cluster-cluster collisions

multiplicity distribution w/o clusters

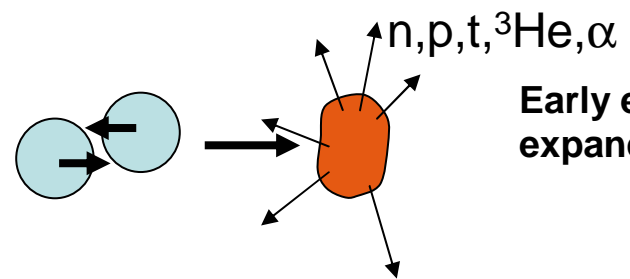


with clusters and cluster-cluster collisions

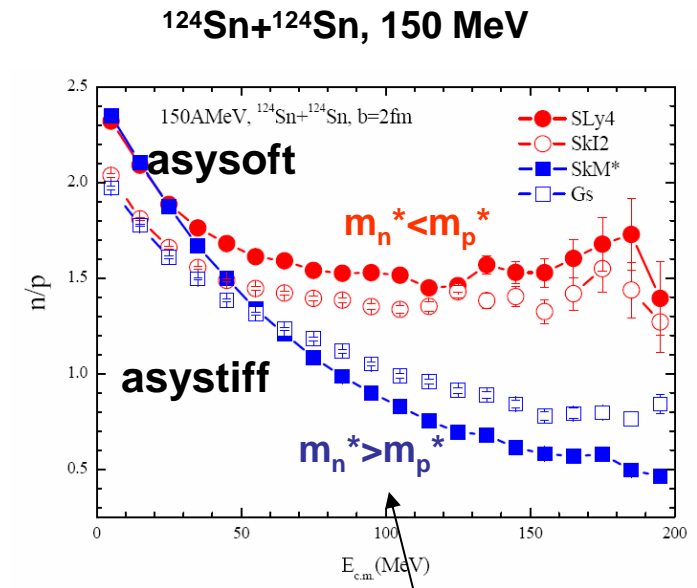
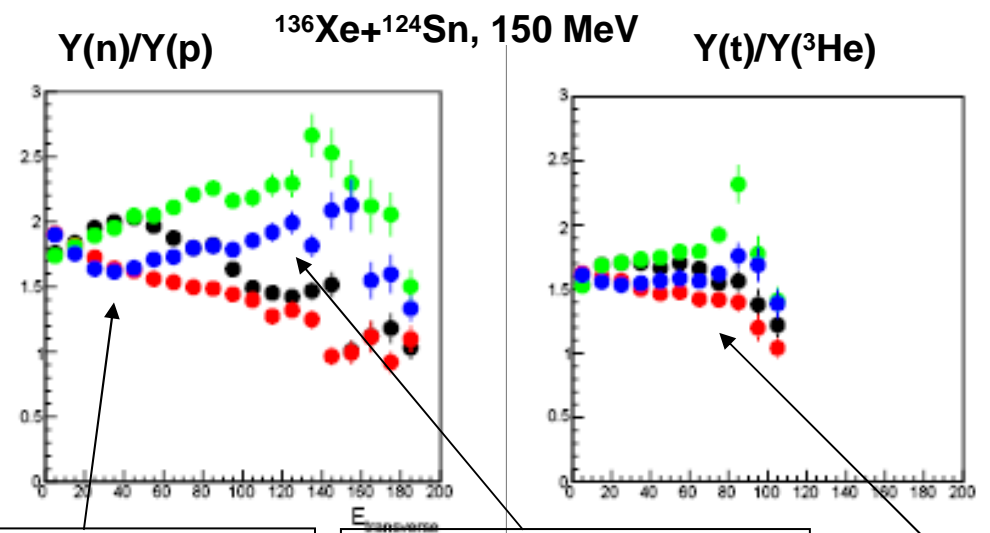


(A. Ono, this meeting)

Example: Ratios of emitted pre-equilibrium particles in central collisions



Early emitted Light Clusters reflect difference in potentials in expanded source, e.g. ratio $Y(n)/Y(p)$.



Y. Zhang, M.B.Tsang, et al., PLB 732, 186 (2014)

Asy-EOS dominates for slow particles; asysoft has larger repulsion at lower densities

Effective mass dominates for fast particles; smaller eff. mass favors emission

Effect also exists for light clusters (easier to measure) but somewhat reduced

similar findings for Sn+Sn collisions (MSU)

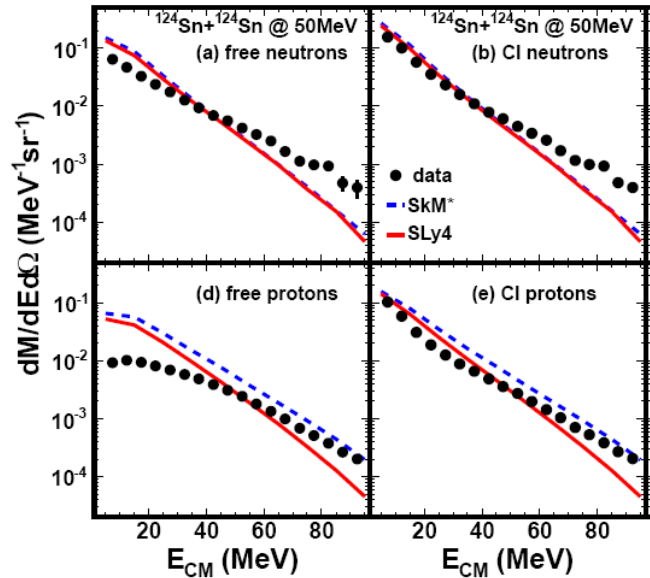
- son: asysoft, $m_n^* > m_p^*$
- stn: asystiff, $m_n^* > m_p^*$
- sop: asysoft, $m_n^* < m_p^*$
- stp: asystiff, $m_n^* < m_p^*$

H.H. Wolter, et al., EPJ Web of Conf. &&, 03097 (2014)

role of clusters?

Comparison with data: problem of light cluster description in transport approaches

free and coalescence invariant (CI) spectra



CI spectra agree better with experiment: A poor man's substitute for not treating light clusters properly in the simulation

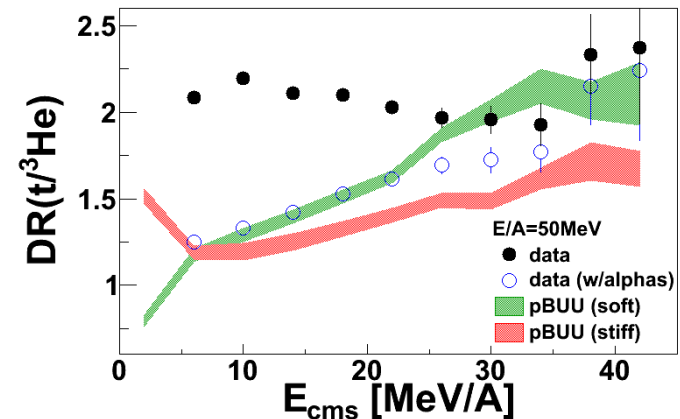
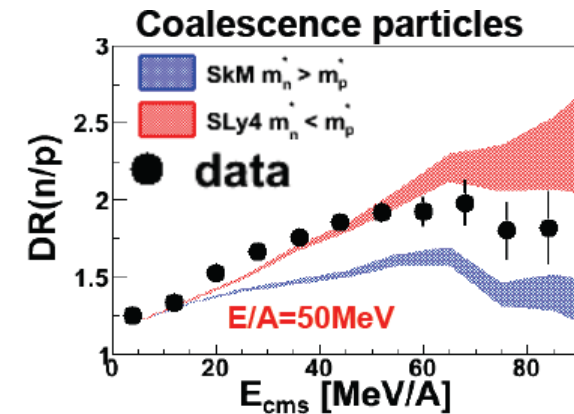
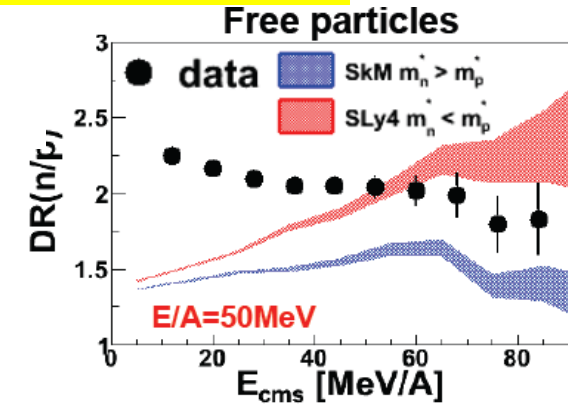
.. or with pBUU calc, if exp. α -particles are counted as t and ^3He ,
→ importance of dynamical LC treatment

W.Lynch, INPC, Florence, 2013:
 $^{124,112}\text{Sn} + ^{124,112}\text{Sn}$, 50 A MeV;
Z. Chajecki, NuSYM13

Double Ratios

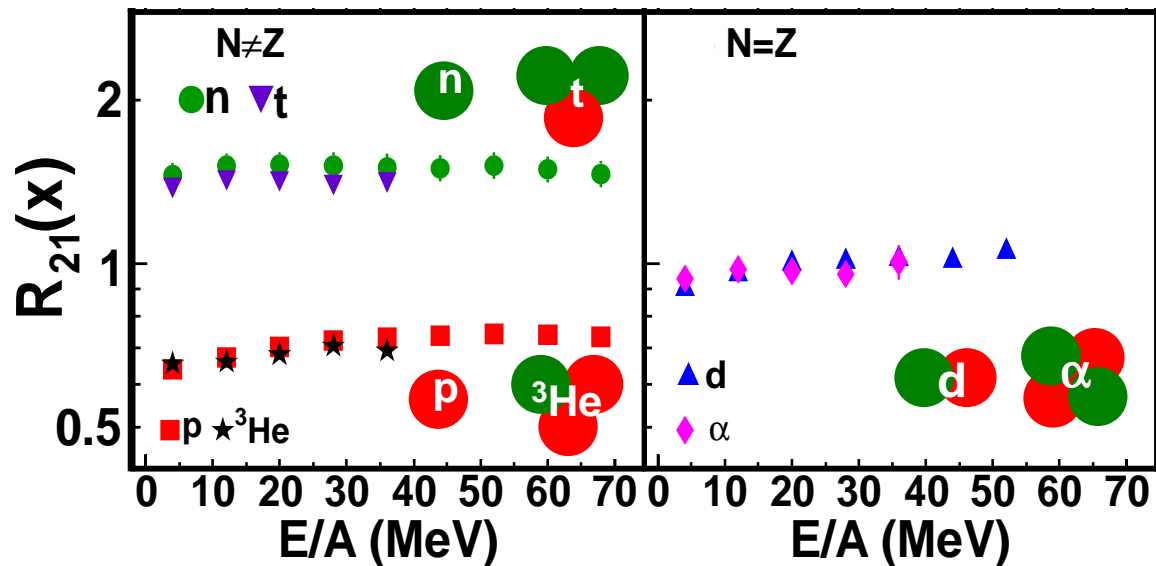
$$\frac{^{124}\text{Sn} + ^{124}\text{Sn}}{^{112}\text{Sn} + ^{112}\text{Sn}}$$

agree only for CI spectra



Y.X. Zhang, M.B.Tsang, et al., PLB 732, 186 (2014)
D.D.S.Coupland, arXiv 1406.4546

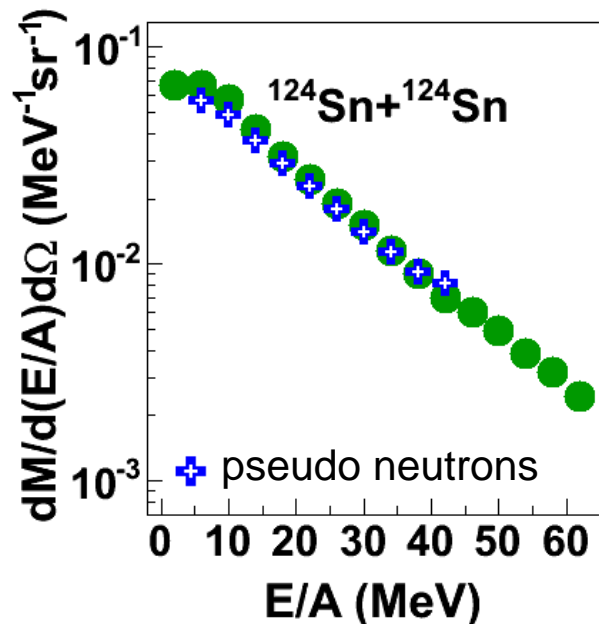
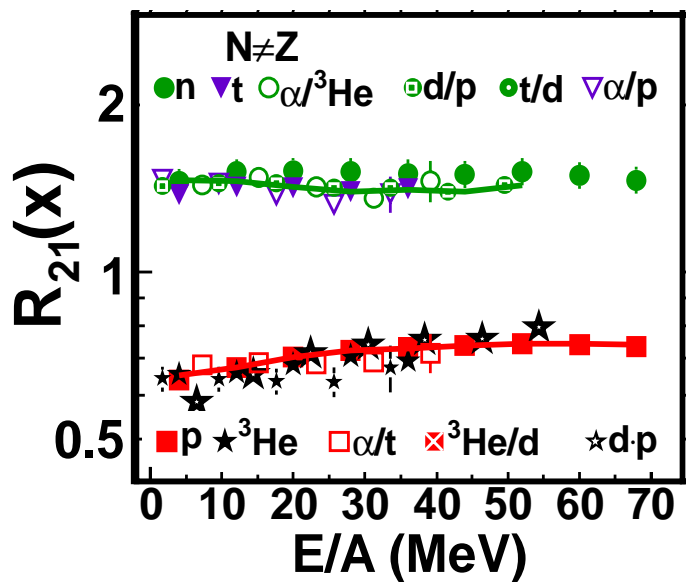
„Chemical potential scaling“ Z.Chajecski, et al., arXiv 1402.5216



$$R_{21}(N, Z) = \frac{dM_2(N, Z)}{dM_1(N, Z)}$$

expect from chemical potential dependence without correlation contributions

$$R_{21}(N, Z) = \exp\left[\frac{(N\Delta\mu_n + Z\Delta\mu_p)}{T}\right]$$

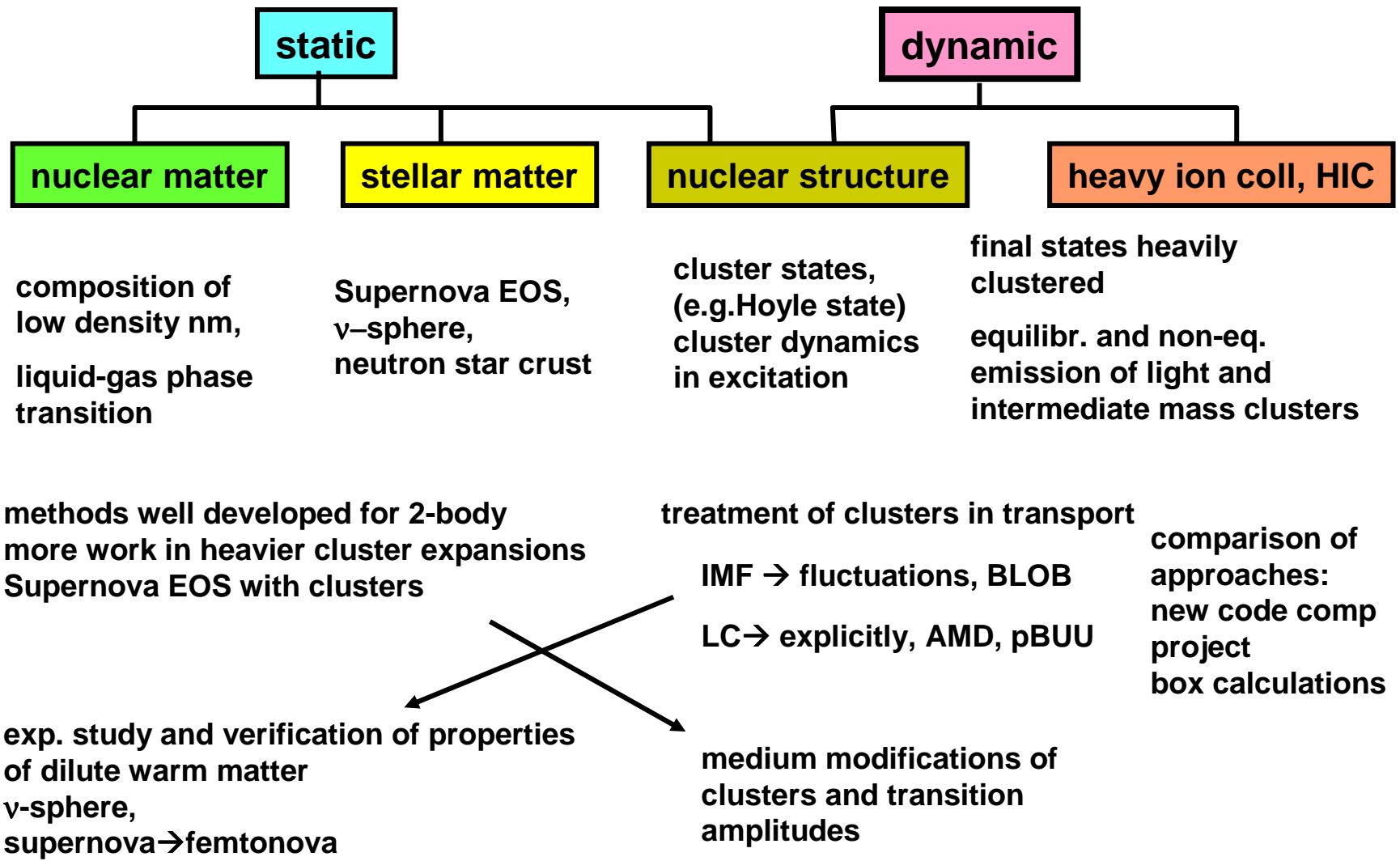


$$\Rightarrow Y(n) = \frac{Y(t)}{Y(^3\text{He})} Y(p)$$

“Pseudo neutron yields”

correlation effects seem to be small for light cluster yields (?)

Cluster Aspects in Nuclear Systems



Clustering enters in many aspects of the investigation of nuclear matter, and esp. the symmetry energy. Here only possible to sketch many ideas, wait for more in-depth talks in this conference!

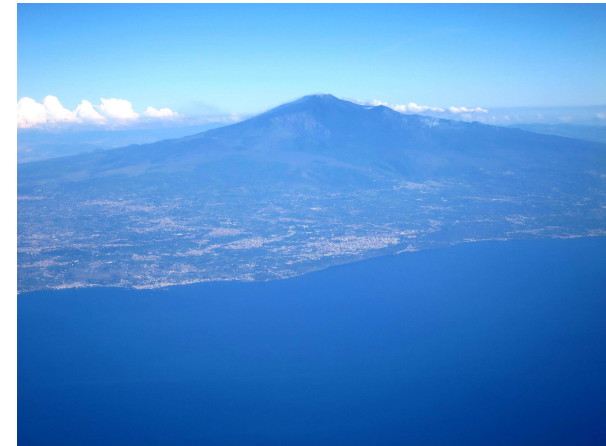
I would like to thank the groups, which whom I have collaborated

:

Catania-Smith connection: M. Colonna, Massimo Di Toro, Enzo Greco, M. Zielinska-Pfabe, and many others

Rostock-Wroclaw-GSI connection: Stefan Typel, Gerd Röpke, David Blaschke, Thomas Klähn, the Texas A&M group,...

MSU-China connection: Pawel Danielewicz, Betty Tsang, W. Lynch, L.W. Chen, BaoAn Li, YingXun Zhang, Jun Xu, ...



...and you for your attention