NUSYM15, Krakow, 30. 6. 2015

Quantum statistical approach to fewnucleon correlations in nuclear systems

Gerd Röpke, Rostock



Quantum statistics

- Symmetry energy: well defined? (Phantasm reality)
- System in equilibrium: temperature T, volume Ω, particle numbers N_c (conserved)
 Thermodynamic potential: free energy F(T, Ω, N_c)
 Internal energy U(T, Ω, N_c)
- Nuclear systems, N_c : neutrons n_n , protons n_p , electrons n_e , ...
- Nuclear structure T=0, astrophysics, heavy ion reactions (HIC): finite T
- Interaction: strong Coulomb
 Separation of the Coulomb part: Interaction energy and structure?

Known results for the EOS

- Low-density limit: ideal quantum gas
- Second virial coefficient: Beth—Uhlenbeck
- Nuclear statistical equilibrium (NSE)
- Cluster-virial expansion

Alpha-particle fraction in the low-density limit

symmetric matter, T=2, 4, 8 MeV



C.J.Horowitz, A.Schwenk, Nucl. Phys. A 776, 55 (2006)

Known results for the EOS

- Low-density limit: ideal quantum gas
- Second virial coefficient: Beth—Uhlenbeck
- Nuclear statistical equilibrium (NSE)
- Cluster-virial expansion
- Saturation density
- Skyrme, relativistic mean-field: Quasiparticles
- Density functional theory

Quasiparticle picture: RMF and DBHF



C. Fuchs, H.H. Wolter, Eur. Phys. J. A 30, 5 (2006)

DBHF at low densities



Difference between **DBHF** calculation and low-density **RMF** fit (square symbols). The corrections (solid lines) are drawn for symmetric nuclear matter (left panel) and pure neutron matter (right panel)

J. Margeron, E. van Dalen, C. Fuchs, Phys. Rev. C 76, 034309 (2007)

Quasiparticle approximation for nuclear matter Equation of state for symmetric matter

10NLo NLoð DBHF DD $D^{2}C$ KVR KVOR DD-F E_0 [MeV] But: cluster -10 formation Incorrect low-density -20^L 0.3 0.2 limit 0.1n [fm⁻³] Klaehn et al., PRC 2006

Symmetry energy and phase transition











Many-particle theory

• Dyson equation and self energy (homogeneous system)

$$G(1, iz_{\nu}) = \frac{1}{iz_{\nu} - E(1) - \Sigma(1, iz_{\nu})}$$

• Evaluation of $\Sigma(1, iz_{\nu})$: perturbation expansion, diagram representation

 $A(1,\omega) = \frac{2 \text{Im } \Sigma(1,\omega+i0)}{[\omega - E(1) - \text{Re } \Sigma(1,\omega)]^2 + [\text{Im } \Sigma(1,\omega+i0)]^2}$ approximation for \longrightarrow approximation for equilibrium correlation functions

alternatively: simulations, path integral methods

Composition of dense nuclear matter

$$n_p(T,\mu_p,\mu_n) = \frac{1}{V} \sum_{A,\nu,K} Z_A f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

$$n_n(T,\mu_p,\mu_n) = \frac{1}{V} \sum_{A,\nu,K} (A - Z_A) f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

mass number A
charge
$$Z_A$$

energy $E_{A,v,K}$
v: internal quantum number $f_{A(z)} = \frac{1}{\exp(z/T) - (-1)^A}$

- Inclusion of excited states and continuum correlations, correct virial expansions
- Medium effects: correct behavior near saturation self-energy and Pauli blocking shifts of binding energies, Coulomb corrections due to screening (Wigner-Seitz,Debye)
- Bose-Einstein condensation, phase instabilities

Few-particle Schrödinger equation in a dense medium

4-particle Schrödinger equation with medium effects

$$\left(\left[E^{HF}(p_1) + E^{HF}(p_2) + E^{HF}(p_3) + E^{HF}(p_4) \right] \right) \Psi_{n,P}(p_1, p_2, p_3, p_4)$$

$$+ \sum_{p_1', p_2'} (1 - f_{p_1} - f_{p_2}) V(p_1, p_2; p_1', p_2') \Psi_{n,P}(p_1', p_2', p_3, p_4)$$

$$+ \left\{ permutations \right\}$$

$$= E_{n,P} \Psi_{n,P}(p_1, p_2, p_3, p_4)$$

Pauli blocking – phase space occupation



cluster wave function (deuteron, alpha,...) in momentum space

P - center of mass momentum

The Fermi sphere is forbidden, deformation of the cluster wave function in dependence on the c.o.m. momentum *P*

momentum space

The deformation is maximal at P = 0. It leads to the weakening of the interaction (disintegration of the bound state).

Shift of the deuteron bound state energy

Dependence on nucleon density, various temperatures, zero center of mass momentum



G.R., NP A 867, 66 (2011)

Shift of the deuteron bound state energy

Dependence on center of mass momentum, various densities, T=10 MeV



G.R., NP A 867, 66 (2011)

Internal energy per nucleon



S. Typel et al., PRC 81, 015803 (2010)

Internal energy per nucleon



EOS for symmetric matter - low density region?

Clustering phenomena in nuclear matter below the saturation density



•FIG. 8. Energy curves of DFSs due to a and 16O clustering in •the symmetric nuclear matter by the use of the BB sB4d force. The •density of matter is normalized by the saturation density of the •uniform matter with the Fermi sphere, r0=0.206 fm-3. The presentation •of the curves is similar to that in Fig. 4.

Hiroki Takemoto et al., PR C **69**, 035802 (2004)

Stefano Gandolfi's talk

EOS of symmetric nuclear matter using Argonne AV6' (no three-body). Low density VERY PRELIMINARY!!!



Gandolfi, Lovato, Carlson, Schmidt, PRC (2014)

Single nucleon distribution function

Dependence on temperature



saturation density

Alm et al., PRC 53, 2181 (1996)

Pauli blocking, correlated medium

In-medium Schroedinger equation

$$[E_{\tau_1}(\mathbf{p}_1; T, \mu_n, \mu_p) + \dots + E_{\tau_A}(\mathbf{p}_A; T, \mu_n, \mu_p) - E_{A\nu}(\mathbf{P}; T, \mu_n, \mu_p)]\psi_{A\nu\mathbf{P}}(1 \dots A) + \sum_{1' \dots A'} \sum_{i < j} [1 - n(i; T, \mu_n, \mu_p) - n(j; T, \mu_n, \mu_p)]V(ij, i'j') \prod_{k \neq i, j} \delta_{kk'}\psi_{A\nu\mathbf{P}}(1' \dots i' \dots j' \dots A') = 0$$

effective occupation numbers

$$n(1) = f_{1,\tau_1}(1) + \sum_{B=2}^{\infty} \sum_{\bar{\nu},\bar{\mathbf{P}}} \sum_{2...B} B f_B \left(E_{B,\bar{\nu}}(\bar{\mathbf{P}};T,\mu_n,\mu_p) \right) |\psi_{B\bar{\nu}\bar{\mathbf{P}}}(1...B)|^2$$

effective Fermi distribution

 $n(1; T, \mu_n, \mu_p) \approx f_{1,\tau_1}(1; T_{\text{eff}}, \mu_n^{\text{eff}}, \mu_p^{\text{eff}}) \qquad \begin{array}{l} \text{blocking by all nucleons} \\ n(1; T, \mu_n, \mu_p) \approx \tilde{f}_{1,\tau_1}(1; T_{\text{eff}}, n_B, Y_p) \\ \end{array}$ effective temperature $T_{\text{eff}} \approx 5.5 \text{ MeV} + 0.5 T + 60 n_B \text{ MeV fm}^3$

G. Roepke, arXiv: 1411.4593, submitted (PRC)

Shift of Binding Energies of Light Clusters







Two-particle correlations



M. Schmidt, G.R., H. Schulz Ann. Phys. **202**, 57 (1990)

FIG. 7. The composition of nuclear matter as a function of the density n for given temperature T = 10 MeV. The solid and dashed lines show the results of the generalized and classical Beth-Uhlenbeck approach, respectively. Note the distinct behavior of n_{free} and n_{corr} predicted by the two approaches in the low and high density limit!



deuteron bound state -2.2 MeV

G. Roepke, J. Phys.: Conf. Series 569, 012031 (2014).

EOS: continuum contributions

Partial density of channel A,c at P (for instance, ${}^{3}S_{1} = d$):

$$z_{A,c}^{\text{part}}(\mathbf{P}; T, \mu_n, \mu_p) = e^{(N\mu_n + Z\mu_p)/T} \left\{ \sum_{\nu_c}^{\text{bound}} g_{A,\nu_c} \ e^{-E_{A,\nu_c}(\mathbf{P})/T} \ \Theta \left[-E_{A,\nu_c}(\mathbf{P}) + E_{A,c}^{\text{cont}}(\mathbf{P}) \right] + z_{A,c}^{\text{cont}}(\mathbf{P}) \right\}$$

separation: bound state part – continuum part ?

$$z_{c}^{\text{part}}(\mathbf{P};T,n_{B},Y_{p}) = e^{[N\mu_{n}+Z\mu_{p}-NE_{n}(\mathbf{P}/A;T,n_{B},Y_{p})-ZE_{p}(\mathbf{P}/A;T,n_{B},Y_{p})]/T} \times g_{c} \left\{ \left[e^{-E_{c}^{\text{intr}}(\mathbf{P};T,n_{B},Y_{p})/T} - 1 \right] \Theta \left[-E_{c}^{\text{intr}}(\mathbf{P};T,n_{B},Y_{p}) \right] + v_{c}(\mathbf{P};T,n_{B},Y_{p}) \right\}$$

parametrization (d – like):

$$v_c(\mathbf{P} = 0; T, n_B, Y_p) \approx \left[1.24 + \left(\frac{1}{v_{T_I=0}(T)} - 1.24\right) e^{\gamma_c n_B/T}\right]^{-1}$$

 $v_d^0(T) = v_{T_I=0}^0(T) \approx 0.30857 + 0.65327 \ e^{-0.102424 \ T/\text{MeV}}$

G. Roepke, arXiv: 1411.4593, submitted (PRC)

Symmetric matter: free energy per nucleon



Dashed lines: no continuum correlations

Internal symmetry energy



Typel et al., PRC 81, 015803 (2010)



Symmetry Energy



Scaled internal symmetry energy as a function of the scaled total density. MDI: Chen et al., QS: quantum statistical, Exp: experiment at TAMU

J.Natowitz et al. PRL, May 2010

Symmetry energy: low density limit



K. Hagel et al., Eur. Phys. J. A (2014) 50: 39

Summary

• The symmetry energy at subsaturation density is strongly depending on temperature T.

• The low-density limit is described by the virial expansion. High values (> 7 MeV) at low temperatures

• The influence of continuum correlations (clusters) at increasing densities requires detailed investigations.

- The blocking of bound states is modified because of correlations in the medium (α matter).

- Continuum correlations contribute to the symmetry energy (density dependent virial coefficients).

• relevant for HIC (freeze-out, transport theory) and astrophysics (supernova explosions)

Thanks

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to you

for attention

D.G.
Symmetry energy, comparison experiment with theories



J.Natowitz et al., PRL 2010

Symmetry energy



Symmetry energy

Heavy-ion collisions, spectra of emitted clusters, temperature (3 - 10 MeV), free energy



Symmetric matter: chemical potential

QS compared with RMF (thin) and NSE (dotted)



G. Roepke, arXiv: 1411.4593, submitted to PRC Insert: no continuum correlations (thin)

Different approximations

• Expansion for small Im $\Sigma(1, \omega + i\eta)$

$$A(1,\omega) \approx \frac{2\pi\delta(\omega - E^{\text{quasi}}(1))}{1 - \frac{d}{dz}\text{Re }\Sigma(1,z)|_{z=E^{\text{quasi}}-\mu_1}} -2\text{Im }\Sigma(1,\omega + i\eta)\frac{d}{d\omega}\frac{P}{\omega + \mu_1 - E^{\text{quasi}}(1)}$$

quasiparticle energy $E^{\text{quasi}}(1) = E(1) + \text{Re} \left[\Sigma(1, \omega) \right]_{\omega = E^{\text{quasi}}}$

• chemical picture: bound states $\hat{=}$ new species



Different approximations

low density limit:

$$G_{2}^{L}(12, 1'2', i\lambda) = \sum_{n\mathbf{P}} \Psi_{n\mathbf{P}}(12) \frac{1}{i\omega_{\lambda} - E_{n\mathbf{P}}} \Psi_{n\mathbf{P}}^{*}(12)$$
$$\sum_{\mathbf{D}} = \begin{bmatrix} \mathbf{T}_{2}^{L} \end{bmatrix}$$

$$n(\beta,\mu) = \sum_{1} f_1(E^{\text{quasi}}(1)) + \sum_{2,n\mathbf{P}}^{\text{bound}} g_{12}(E_{n\mathbf{P}}) + \sum_{2,n\mathbf{P}} \int_0^\infty \mathrm{d}k \ \delta_{\mathbf{k},\mathbf{p}_1-\mathbf{p}_2} g_{12}(E^{\text{quasi}}(1) + E^{\text{quasi}}(2)) 2 \sin^2 \delta_n(k) \frac{1}{\pi} \frac{\mathrm{d}}{\mathrm{d}k} \delta_n(k)$$

• generalized Beth-Uhlenbeck formula correct low density/low temperature limit: mixture of free particles and bound clusters

M. Schmidt et al., Ann. Phys. (NY) 202, 57 (1990)

Effective wave equation for the deuteron in matter

In-medium two-particle wave equation in mean-field approximation $\left(\frac{p_{1}^{2}}{2m_{1}} + \Delta_{1} + \frac{p_{2}^{2}}{2m_{2}} + \Delta_{2}\right)\Psi_{d,P}(p_{1},p_{2}) + \sum_{p_{1}',p_{2}'}(1 - f_{p_{1}} - f_{p_{2}})V(p_{1},p_{2};p_{1}',p_{2}')\Psi_{d,P}(p_{1}',p_{2}')$

Add self-energy

Pauli-blocking

$$= E_{d,P} \Psi_{d,P}(p_1,p_2)$$

Thouless criterion $E_d(T,\mu) = 2\mu$

Fermi distribution function

$$f_p = \left[e^{(p^2/2m - \mu)/k_B T} + 1 \right]^{-1}$$

BEC-BCS crossover: Alm et al.,1993

EOS: continuum contributions

Partial density of channel A,c at P (for instance, ${}^{3}S_{1} = d$):

$$z_{A,c}^{\text{part}}(\mathbf{P}; T, \mu_n, \mu_p) = e^{(N\mu_n + Z\mu_p)/T} \left\{ \sum_{\nu_c}^{\text{bound}} g_{A,\nu_c} \ e^{-E_{A,\nu_c}(\mathbf{P})/T} \ \Theta \left[-E_{A,\nu_c}(\mathbf{P}) + E_{A,c}^{\text{cont}}(\mathbf{P}) \right] + z_{A,c}^{\text{cont}}(\mathbf{P}) \right\}$$

separation: bound state part – continuum part ?

$$z_{c}^{\text{part}}(\mathbf{P};T,n_{B},Y_{p}) = e^{[N\mu_{n}+Z\mu_{p}-NE_{n}(\mathbf{P}/A;T,n_{B},Y_{p})-ZE_{p}(\mathbf{P}/A;T,n_{B},Y_{p})]/T} \times g_{c} \left\{ \left[e^{-E_{c}^{\text{intr}}(\mathbf{P};T,n_{B},Y_{p})/T} - 1 \right] \Theta \left[-E_{c}^{\text{intr}}(\mathbf{P};T,n_{B},Y_{p}) \right] + v_{c}(\mathbf{P};T,n_{B},Y_{p}) \right\}$$

parametrisierung (d – like):

$$v_c(\mathbf{P}=0;T,n_B,Y_p) \approx \left[1.24 + \left(\frac{1}{v_{T_I=0}(T)} - 1.24\right)e^{\gamma_c n_B/T}\right]^{-1}$$

 $v_d^0(T) = v_{T_I=0}^0(T) \approx 0.30857 + 0.65327 \ e^{-0.102424 \ T/\text{MeV}}$

 $\gamma_d(\mathbf{P}, T, n_B, Y_p) = 1873.2 \text{ MeV fm}^3 \exp\left[-P^2 \text{ fm}^2/(1.8463 + 0.1617 T \text{ MeV}^{-1} + 0.17 P^2 \text{ fm}^2)\right]$

G. Roepke, arXiv: 1411.4593, submitted (PRC)

Symmetric nuclear matter: Phase diagram



Nuclear matter phase diagram



Cluster - mean field approximation

Cluster (A) interacting with a distribution of clusters (B) in the medium, fully antisymmetrized

$$\sum_{1'\dots A'} \{H_A^0(1\dots A, 1'\dots A') + \sum_i \Delta_i^{A,mf} \delta_{k,k'} + \frac{1}{2} \sum_{i,j} \Delta V_{ij}^{A,mf} \delta_{l,l'} - E_{AvP} \delta_{k,k'} \} \psi_{AvP}(1'\dots A') = 0$$

self-energy

$$\Delta_1^{A,mf}(1) = \sum_2 V(12,12)_{ex} f^*(2) + \sum_{BvP} \sum_{2...B'} f_B(E_{BvP}) \sum_i V_{1i}(1i,1'i') \psi_{BvP}^*(1...B) \psi_{BvP}(1'...B')$$

effective interaction

$$\Delta V_{12}^{A,mf} = -\frac{1}{2} [f^*(1) + f^*(2)] V(12,1'2') - \sum_{B\nu P} \sum_{2^* \dots B^*} f_B(E_{B\nu P}) \sum_i V_{1i} \psi_{B\nu P}^*(22^* \dots B^*) \psi_{B\nu P}(2'2" \dots B")$$

phase space occupation $f^*(1) = f_1(1) + \sum_{B \lor P} \sum_{2...B} f_B(E_{B \lor P}) |\psi_{B \lor P}(1...B)|^2$

Single nucleon distribution function

Dependence on density



T = 10 MeV

Alm et al., PRC 53, 2181 (1996)

Different approximations

Ideal Fermi gas:

protons, neutrons, (electrons, neutrinos,...)

bound state formation

Nuclear statistical equilibrium: ideal mixture of all bound states (clusters:) chemical equilibrium

continuum contribution

Second virial coefficient: account of continuum contribution, scattering phase shifts, Beth-Uhl.E.

chemical & physical picture

Cluster virial approach:

all bound states (clusters) scattering phase shifts of all pairs

medium effects

Quasiparticle quantum liquid: mean-field approximation Skyrme, Gogny, RMF

Chemical equilibrium with quasiparticle clusters: self-energy and Pauli blocking

Generalized Beth-Uhlenbeck formula:

medium modified binding energies, medium modified scattering phase shifts

Correlated medium

Symmetric matter: composition



G. Roepke, arXiv: 1411.4593, submitted to PRC

Free symmetry energy



symmetry entropy

Internal symmetry energy

R. Wada et al., Phys. Rev. C 85, 064618 (2012).

Chemical potential



preliminary

Cluster virial expansion for nuclear matter within a quasiparticle approach

Generalized Beth-Uhlenbeck approach

$$n_1^{\rm qu}(T,\mu_p,\mu_n) = \sum_{A,Z,\nu} \frac{A}{\Omega} \sum_{\substack{\vec{P} \\ P > P_{\rm Mott}}} f_A(E_{A,Z,\nu}(\vec{P};T,\mu_p,\mu_n),\mu_{A,Z,\nu})$$

$$n_{2}^{qu}(T,\mu_{p},\mu_{n}) = \sum_{A,Z,\nu} \sum_{A',Z',\nu'} \frac{A+A'}{\Omega} \sum_{\vec{p}} \sum_{c} g_{c} \frac{1+\delta_{A,Z,\nu;A',Z',\nu'}}{2\pi} \times \int_{0}^{\infty} dE f_{A+A'} \left(E_{c}(\vec{P};T,\mu_{p},\mu_{n}) + E,\mu_{A,Z} + \mu_{A',Z'} \right) 2 \sin^{2}(\delta_{c}) \frac{d\delta_{c}}{dE}$$

Avoid double counting



Generating functional



G.R., N. Bastian, D. Blaschke, T. Klaehn, S. Typel, H. Wolter, NPA 897, 70 (2013)

Chemical picture and medium corrections

Nuclear matter at given temperature T, baryon density n_B , proton fraction (asymmetry) $Y_e = n_p/n_B$: equation of state

- Iow-density limit: ideal mixture of reacting components: Nuclear statistical equilibrium (NSE)
- interactions: virial expansion (cluster virial expansion)
- higher densities: quasiparticle concept, medium modification of components (cluster mean-field approximation)
 - nucleons as quasiparticles:
 - Skyrme, relativistic mean-field (RMF), Dirac Brueckner Hartree-Fock
 - light elements (d, t, h, α) as quasiparticles: shift of energy (self-energy, Pauli blocking), Mott effect.
 - excluded volume
 - quantum statistical approach (QS): E_{A,Z}(p;T,n_B,Y_e)

Nuclear matter phase diagram





Free energy per nucleon



Constrained THSR calculations as function of the c.o.m. width B?

Chemical potential of symmetric matter



Light Cluster Abundances



S. Typel et al., PRC **81**, 015803 (2010)

Cluster yields in HIC

PRL 108, 062702 (2012)

PHYSICAL REVIEW LETTERS

week ending 10 FEBRUARY 2012

Experimental Determination of In-Medium Cluster Binding Energies and Mott Points in Nuclear Matter

K. Hagel,¹ R. Wada,^{2,1} L. Qin,¹ J. B. Natowitz,¹ S. Shlomo,¹ A. Bonasera,^{1,3} G. Röpke,⁴ S. Typel,⁵ Z. Chen,² M. Huang,² J. Wang,² H. Zheng,¹ S. Kowalski,⁶ C. Bottosso,¹ M. Barbui,¹ M. R. D. Rodrigues,¹ K. Schmidt,¹ D. Fabris,⁷ M. Lunardon,⁷ S. Moretto,⁷ G. Nebbia,⁷ S. Pesente,⁷ V. Rizzi,⁷ G. Viesti,⁷ M. Cinausero,⁸ G. Prete,⁸ T. Keutgen,⁹ Y. El Masri,⁹ and Z. Majka¹⁰



Mott points from cluster yields



FIG. 3. Comparison of experimentally derived Mott point densities and temperatures with theoretical values. Symbols represent the experimental data. Estimated errors on the temperatures are 10% and on the densities 20%. Lines show polynomial fits to the Mott points presented in Ref. [1].

K. Hagel et al., PRL **108**, 062702 (2012)

Different approximations

Ideal Fermi gas:

protons, neutrons, (electrons, neutrinos,...)

bound state formation

Nuclear statistical equilibrium: ideal mixture of all bound states (clusters:) chemical equilibrium

continuum contribution

Second virial coefficient: account of continuum contribution, scattering phase shifts, Beth-Uhl.E.

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Generalized Beth-Uhlenbeck formula:

medium modified binding energies, medium modified scattering phase shifts

Correlated medium

Cluster decomposition of the self-energy



T-matrices: bound states, scattering states Including clusters like new components chemical picture, mass action law, nuclear statistical equilibrium (NSE)

Beth-Uhlenbeck formula

rigorous results at low density: virial expansion

Beth-Uhlenbeck formula

$$n(T,\mu) = \frac{1}{V} \sum_{p} e^{-(E(p)-\mu)/k_{B}T} + \frac{1}{V} \sum_{nP} e^{-(E_{nP}-2\mu)/k_{B}T} + \frac{1}{V} \sum_{\alpha P} \int_{0}^{\infty} \frac{dE}{2\pi} e^{-(E+P^{2}/4m-2\mu)/k_{B}T} \frac{d}{dE} \delta_{\alpha}(E) + \dots$$

 $\delta_{\alpha}(E)$: scattering phase shifts, channel α

Quantum condensates in nuclei?

Lot of semantics – my position

- Pairing is well accepted.
- Quartetting is not very well-known and simple.
- The main point is the formation of clusters (correlations) in lowdensity matter.
- We are interested in an efficient description (optimal wave function) for the cluster state.
- The center of mass motion has to be considered as new (collective) degree of freedom.

Four-nucleon energies at finite density

Solution of the in-medium wave equation, T = 0



Pauli blocking and Mott effect

Two different fermions (a,b: proton, neutron) form a bound state (c: deuteron).

$$c_{q} = \sum_{p} F(q,p)a_{p}b_{q-p}$$
Is the bound state a boson? Commutator relation
$$\begin{bmatrix} c_{q}, c_{q'}^{+} \end{bmatrix}_{-} = \sum_{p,p'} F(q,p)F^{*}(q',p')\begin{bmatrix} a_{p}b_{q-p}, b_{q'-p'}^{+}a_{p'}^{+} \end{bmatrix}_{-}$$

$$\frac{a_{p}b_{q-p}b_{q'-p'}^{+}a_{p'}^{+} + a_{p}b_{q'-p'}^{+}a_{p'}^{+} - a_{p}b_{q'-p'}^{+}b_{q-p}a_{p'}^{+} - b_{q'-p'}^{+}a_{p}b_{q-p}a_{p'}^{+} + b_{q'-p'}^{+}a_{p}a_{p}^{+}b_{q-p} - b_{q'-p'}^{+}a_{p}a_{p}^{+}b_{q-p} - b_{q'-p'}^{+}a_{p}b_{q-p}a_{p}^{+}a_{p}b_{q-p}a_{p}^{+} + b_{q'-p'}^{+}a_{p}a_{p}^{+}b_{q-p} - b_{q'-p'}^{+}a_{p}a_{p}^{+}b_{q-p} - b_{q'-p'}^{+}a_{p}b_{q-p}a_{p}^{+}a_{p}b_{q-p}a_{p}^{+}a_{p}^{+}b_{q-p}^{-}a_{p}^{+}a_{p}^{+}b_{q-p}^{-}b_{q'-p'}^{+}a_{p}b_{q-p}a_{p}^{+}b_{q-p} - b_{q'-p'}^{+}a_{p}a_{p}b_{q-p}a_{p}^{+}a_{p}b_{q-p}a_{p}^{+}b_{q-p}^{-}a_{p}^{+}a_{p}b_{q-p} - b_{q'-p'}^{+}a_{p}b_{q-p}b_{q-p}b_{q-p}b_{p,p'}a_{p}b_{q-p}a_{p}^{+}b_{q-p}b_{q-p}b_{p,p'} = \left(b_{p,p'} - a_{p'}^{+}a_{p}\right)\delta_{q-p,q'-p'} - b_{q'-p'}^{+}b_{q'-p'}b_{q-p}\delta_{p,p'}\right]$$

$$\begin{bmatrix} c_{q}, c_{q'}^{+} \end{bmatrix}_{-} = \sum_{p,p'} F(q,p)F^{*}(q',p') \left[\left(\delta_{p,p'} - a_{p'}^{+}a_{p}\right) \delta_{q-p,q'-p'} - b_{q'-p'}^{+}b_{q-p}b_{p,p'} \right]$$

$$= \sum_{p} F(q,p)F^{*}(q,p)\delta_{q,q'} - \sum_{p,p'} F(q,p)F^{*}(q',p') \left[\left(a_{p'}^{+}a_{p}\right) \delta_{q-p,q'-p'} + \left(b_{q'-p'}^{+}b_{q-p}\right) \delta_{p,p'} \right]$$
averaging

averaging

$$\left\langle \left[c_{q},c_{q'}^{+}\right]_{-}\right\rangle = \delta_{q,q'}\left[1 - \sum_{p} F(q,p)F^{*}(q,p)\left(\left\langle a_{p}^{+}a_{p}\right\rangle + \left\langle b_{q-p}^{+}b_{q-p}\right\rangle\right)\right]$$

Fermionic substructure: phase space occupation, "excluded volume"

Ideal mixture of reacting nuclides

$$n_p(T,\mu_p,\mu_n) = \frac{1}{V} \sum_{A,\nu,K} Z_A f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

$$n_n(T,\mu_p,\mu_n) = \frac{1}{V} \sum_{A,\nu,K} (A - Z_A) f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

$$f_{A(z)} = \frac{1}{\exp(z/T) - (-1)^A}$$

mass number A, charge Z_A , energy $E_{A,v,K}$, v internal quantum number,

Nuclear Statistical Equilibrium (NSE)

Parametrization

- Single-nucleon quasiparticle energies
 E_τ(p, T, n_B, Y_e)
 (DBHF, Skyrme, RMF,...)
- Bound state energies $E_{A,Z,v}(p, T, n_{B}, Y_e)$

G.R., NP A 867, 66 (2011)

Composition of dense nuclear matter

$$\begin{split} n_{p}(T,\mu_{p},\mu_{n}) &= \frac{1}{V} \sum_{A,\nu,K} Z_{A} f_{A} \{ E_{A,\nu K} - Z_{A} \mu_{p} - (A - Z_{A}) \mu_{n} \} \\ n_{n}(T,\mu_{p},\mu_{n}) &= \frac{1}{V} \sum_{A,\nu,K} (A - Z_{A}) f_{A} \{ E_{A,\nu K} - Z_{A} \mu_{p} - (A - Z_{A}) \mu_{n} \} \\ \max_{\substack{\text{charge } Z_{A}, \\ \text{energy } E_{A,\nu,K}, \\ \text{v: internal quantum number,}} f_{A(z)} &= \frac{1}{\exp(z/T) - (-1)^{A}} \end{split}$$

- Inclusion of excited states and continuum correlations
- Medium effects:

self-energy and Pauli blocking shifts of binding energies, Coulomb corrections due to screening (Wigner-Seitz,Debye)

Supernova explosion



Core-collapse supernovae



Density.

electron fraction, and

temperature profile

of a 15 solar mass supernova at 150 ms after core bounce as function of the radius.

Influence of cluster formation on neutrino emission in the cooling region and on neutrino absorption in the heating region ?

K.Sumiyoshi et al., Astrophys.J. **629**, 922 (2005)
Composition of supernova core



K.Sumiyoshi, G. R., PRC **77**, 055804 (2008)

Mass fraction X of light clusters for a post-bounce supernova core

Composition of supernova core



r [km]

S. Heckel, P. P. Schneider and A. Sedrakian, Light nuclei in supernova envelopes: a quasiparticle gas model Phys. Rev. C **80**, 015805 (2009).

alpha-fraction in symmetric matter



α - α scattering phase shifts



C.J.Horowitz, A.Schwenk, Nucl. Phys. A 776, 55 (2006)

α -n scattering phase shifts



Fig. 2. (Color online.) The phase shifts for elastic neutron-alpha scattering $\delta_{L_J}(E)$ versus laboratory energy *E*. As discussed in the text, the solid lines are from Arndt and Roper [37] and the symbols are from Amos and Karataglidis [38]. For clarity, we do not show the F-waves included in our results for $b_{\alpha n}$.

C.J.Horowitz, A.Schwenk, Nucl. Phys. A 776, 55 (2006)

Correlations in the medium



cluster mean-field approximation

EOS at low densities from HIC



Influence of cluster formation on β equilibrium



Influence of cluster formation on β equilibrium



Liquid-vapor phase transition



blue: no light cluster, green: with light clusters, QS, red: cluster-RMF S. Typel et al., PRC 81, 015803 (2010)

α -matter and cluster-virial expansion

- Bound states may become dominant: Treat it like new species (chemical picture)
- Cluster expansion: Introduce corresponding fugacities (chemical potentials)
- Inclusion of continuum states: scattering phase shifts, Beth-Uhlenbeck formula
- Introduce quasiparticles and avoid double counting; transition to the high-density region

Bound state formation in many-fermion systems

Examples:

Quark-gluon plasma: hadrons; deconfinement phase transition Nucleon systems: nuclei; nuclear matter Electron-ion plasma: atoms; metals Electron-hole plasma: exciton; electron-hole droplets

Properties:

Equation of state, transport coefficient, dynamical structure factor

Theoretical approaches:

Simulations, Path integral, Quantum statistics of many-particle systems:

- Lagrangian/Hamiltonian
- Green functions (Feynman diagrams, partial summations, chemical picture)
- Self-energy, spectral function
- Correlation functions, physical properties

Nucleon-nucleon interaction

QCD? Effective Lagrangians, interaction potentials (PEST)

singlet (nn, pp): a = -23.678 fm, r = 1.726 fm triplet (pn): a = 5.396 fm, r = 2.729 fm, E = -2.225 MeV $k \cot \delta = -\frac{1}{a} + r_0 \frac{k^2}{2}$

Separable interaction

• general form:

$$egin{aligned} V_lpha(p,p') &= \sum\limits_{i,j=1}^N w_{lpha i}(p)\lambda_{lpha ij}w_{lpha j}(p') & ext{uncoupled} \ & ext{and} \ V^{LL'}_lpha(p,p') &= \sum\limits_{i,j=1}^N w^L_{lpha i}(p)\lambda_{lpha ij}w^{L'}_{lpha j}(p') & ext{coupled} \end{aligned}$$

- p, p' in- and outgoing relative momentum
- lpha ... channel
- N ... rank
- $\lambda_{lpha ij}$. coupling parameter
- L, L' orbital angular momentum

Weak interaction - beta equilibrium? Coulomb interaction?

Composition of symmetric nuclear matter

Fraction of correlated matter (virial expansion, Generalized Beth-Uhlenbeck approach, contribution of bound states, of scattering states, phase shifts)

H. Stein et al., Z. Phys. **A351**, 259 (1995)



QS versus NSE: comparison with data

40Ar124Sn K_{alpha}



Energy of α -Matter at T=0



Total energy calculated with the cluster expansion within the HNC/0 (circles) and HNC/4 (solid lines) approximation. Different interaction potentials

F.Carstoiu, S.Misicu, PLB, 2009

Composition of dense nuclear matter

$$n_p(T,\mu_p,\mu_n) = \frac{1}{V} \sum_{A,\nu,K} Z_A f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

$$n_n(T,\mu_p,\mu_n) = \frac{1}{V} \sum_{A,\nu,K} (A - Z_A) f_A \{ E_{A,\nu K} - Z_A \mu_p - (A - Z_A) \mu_n \}$$

mass number A
charge
$$Z_A$$

energy $E_{A,v,K}$
v: internal quantum number
 $f_A(z) = \frac{1}{\exp(z/T) - (-1)^A}$

- Inclusion of excited states and continuum correlations
- Medium effects:

self-energy and Pauli blocking shifts of binding energies, Coulomb corrections due to screening (Wigner-Seitz,Debye)

Bose-Einstein condensation

Deuterons in nuclear matter



T=10 MeV, P: center of mass momentum

Scattering phase shifts in matter



Outline

- Correlations and bound state formation in fermion systems The chemical picture, Green functions, spectral functions
- Correlations in nuclear systems

where it occurs, what do we know: nuclei, stars, HIC; effective interactions

- Many-particle theory: Equation of state, generalized Beth-Uhlenbeck eq. Low-density limit: Mass action law, nuclear statistical equilibrium, virial expansion Near saturation: quasiparticles, dissolution of bound states, correlated mean field
- Applications:

Equation of state in SN explosions, HIC, Hoyle-like states, surface correlations

- Quantum condensates in nuclear matter: Pairing vs. quartetting, BEC-BCS crossover
- Bose condensates in finite systems:

Suppression of condensate, HFB and cluster formation (pairing and quartetting)

Applications

Signatures of clustering

- Nuclear structure: Hoyle state, states near the threshold n-alpha decay
- rms radii: comparatively large, low-density matter, skin?
- Nuclear reactions, alpha decay (preformation at the surface)
- Yields of light elements in HIC, equation of state and symmetry energy
- Excited states

Clustering is not "exotic".

It is a general feature of (low density) nuclear systems.

Mott points from cluster yields

PRL 108, 062702 (2012)

PHYSICAL REVIEW LETTERS

week ending 10 FEBRUARY 2012

Experimental Determination of In-Medium Cluster Binding Energies and Mott Points in Nuclear Matter

K. Hagel,¹ R. Wada,^{2,1} L. Qin,¹ J. B. Natowitz,¹ S. Shlomo,¹ A. Bonasera,^{1,3} G. Röpke,⁴ S. Typel,⁵ Z. Chen,² M. Huang,² J. Wang,² H. Zheng,¹ S. Kowalski,⁶ C. Bottosso,¹ M. Barbui,¹ M. R. D. Rodrigues,¹ K. Schmidt,¹ D. Fabris,⁷ M. Lunardon,⁷ S. Moretto,⁷ G. Nebbia,⁷ S. Pesente,⁷ V. Rizzi,⁷ G. Viesti,⁷ M. Cinausero,⁸ G. Prete,⁸ T. Keutgen,⁹ Y. El Masri,⁹ and Z. Majka¹⁰



FIG. 3. Comparison of experimentally derived Mott point densities and temperatures with theoretical values. Symbols represent the experimental data. Estimated errors on the temperatures are 10% and on the densities 20%. Lines show polynomial fits to the Mott points presented in Ref. [1].

Symmetry energy at medium densities

The Nuclear Matter Symmetry Energy at $0.03 \le \rho/\rho_0 \le 0.2$

R. Wada,^{1,2} K. Hagel,² L. Qin,² J. B. Natowitz,² Y. G. Ma,³ G. Röpke,⁴ S. Shlomo,² A. Bonasera,^{2,5} S.

Typel,⁶ Z. Chen,^{2,1} M. Huang,^{2,1} J. Wang,^{2,1} H. Zheng,² S. Kowalski,⁷ C. Bottosso,² M. Barbui,²

M. R. D. Rodrigues,² K. Schmidt,² D. Fabris,⁸ M. Lunardon,⁸ S. Moretto,⁸ G. Nebbia,⁸ S. Pesente,⁸

V. Rizzi,⁸ G. Viesti,⁸ M. Cinausero,⁹ G. Prete,⁹ T. Keutgen,¹⁰ Y. El Masri,¹⁰ and Z. Majka¹¹



In-medium modification of transport properties of dense matter

- D. Blaschke, G. Röpke, H. Schulz, A.D. Sedrakian, Nuclear in-medium effect on the thermal conductivity and viscosity of neutron star matter
 PL B 338, 111 (1994)
- D. Blaschke, G. Röpke, H. Schulz, A.D. Sedrakian, D. Voskresensky, Nuclear in-medium effects and neutrino emissivity of neutron stars.
 M. N. R. A. S. 273, 596 (1995)

Supernova Crab nebula, 1054 China, PSR 0531+21



M1, the Crab Nebula. Courtesy of NASA/ESA

Nuclear matter equation of state

- Nuclear systems: Quasiparticle approach Brueckner, HFB; Skyrme, Relativistic Mean Field (RMF)
- Account of correlations in warm dense matter: two-particle (deuteron, pairing), four-particle (alpha-like) correlations, light elements
- Low-density regions: Nuclear Statistical Equilibrium (NSE) Hoyle-like states in light expanded nuclei, surface of nuclei, neck emission, alpha matter...
- Quantum statistical approach (n < 0.15 fm⁻³, T < 20 MeV) Equation of state, Beth-Uhlenbeck formula disappearance of clusters at high densities, Pauli blocking
- Experimental signatures Heavy Ion Collisions (HIC), Symmetry energy, SN explosions, ...

Diffusion Monte Carlo EOS calculation



S. Gandolfi, A. Yu. Illarionov, et al., Mon.Not.R.Astron.Soc., 2010